Direct Higgs production at hadron colliders

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Les Houches 2003

- Inclusive cross section
  - Soft-gluon resummation at NNLL
  - Residual theoretical uncertainty
- Transverse momentum distribution
  - An improved resummation formalism
  - NNLL+NLO results

Motivations

- Several constraints indicate $100 < M_H < 200$ GeV
- SM Higgs boson with $M_H \sim 115$ GeV 'found' at LEP?
- Next search at Hadron Colliders: Run II at Tevatron and LHC
- Dominant production mechanism: gg fusion through a heavy-quark loop
- NLO QCD corrections known to be large

A.Djouadi, M.Spira, P.Zerwas (1991)
S.Dawson (1991)

- Estimates of higher order corrections still sizable

M.Krämer, E.Laenen, M.Spira (1998)

$\Downarrow$

NNLO calculation highly demanded

Direct Higgs production at NNLO

Evaluation of NNLO corrections completed

R.Harlander (2000)
S.Catani, D. de Florian, MG (2001)
R.Harlander, W.Kilgore (2001,2002)
C.Anastasiou and K.Melnikov (2002)
V.Ravindran, J.Smith, W.L. van Neerven (2003)

Together with approximated NNLO pdf's available:
MRST2002, Alekhin (kernels from van Neerven, Vogt)

$\Rightarrow$ Consistent calculation at NNLO

Effect wrt NLO: +15% at the LHC
+35% at the Tevatron

The soft + virtual + leading collinear (SVC) contributions provide an excellent approximation of the full result

The hard contribution is about 2% at LHC and 4% at the Tevatron

Why does the soft approximation work so well ?

The hard cross section is convoluted with the parton distributions that are strongly peaked at small x:
$< \hat{s} >=< x_1 x_2 > s \ll s$

$\Rightarrow$The hard cross section is almost always evaluated close to threshold

The $M_H \ll m_{top}$ approximation

For a light Higgs it is possible to use an effective lagrangian in the limit $m_{top} \to \infty$:

$$\mathcal{L}_{eff} = -\frac{1}{4}\left[1 - \frac{\alpha_S}{3\pi}\frac{H}{v}(1+\Delta)\right] Tr G_{\mu\nu}G^{\mu\nu}$$

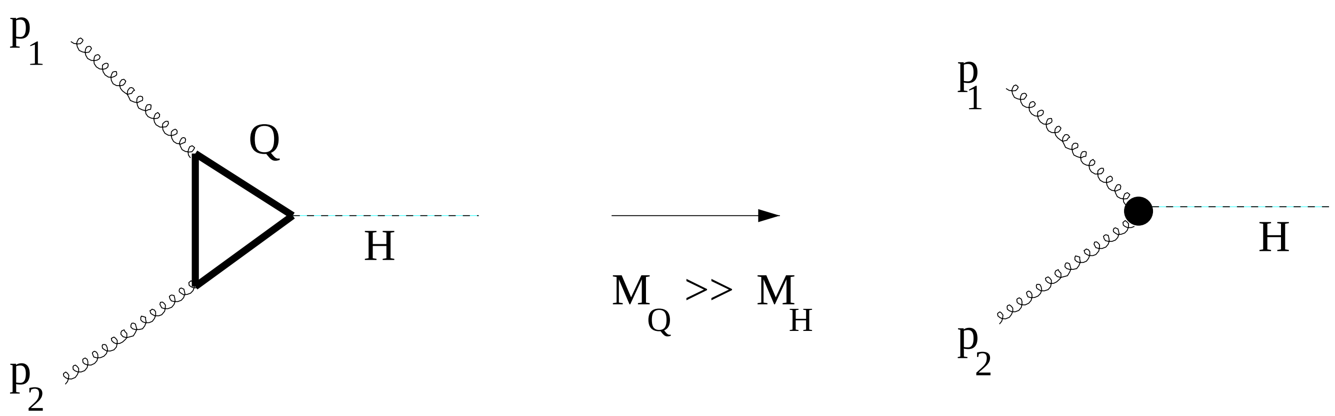
J.Ellis, M.K.Gaillard, D.V.Nanopoulos (1976)
M.Voloshin, V.Zakharov, M.Shifman (1979)

The correction Δ is known up to $\mathcal{O}(\alpha_S^3)$

K.G.Chetyrkin, B.A.Kniehl, M.Steinhauser (1997)

The approximation is correct within 5% if $M_H < 2\, m_{top}$

Effective vertex: one loop less



The approximation works well even for large M_H if full Born result is retained

Soft gluon resummation

S.Catani, D.de Florian, P.Nason, MG (2002)

Inclusive cross section dominated by soft and collinear emission $\Rightarrow$Multiple soft emission beyond NNLO can be important

In N-space the large logarithmic corrections appear as $\alpha_S^n \log^{2n} N$

This large corrections can be resummed to all orders:

$$G_{gg\to H,N}(\alpha_S) = C(\alpha_S)\exp\Big\{\ln N g_1(\beta_0\alpha_S \ln N) + g_2(\beta_0\alpha_S \ln N) + \alpha_S\, g_3(\beta_0\alpha_S \ln N) + \dots\Big\}$$

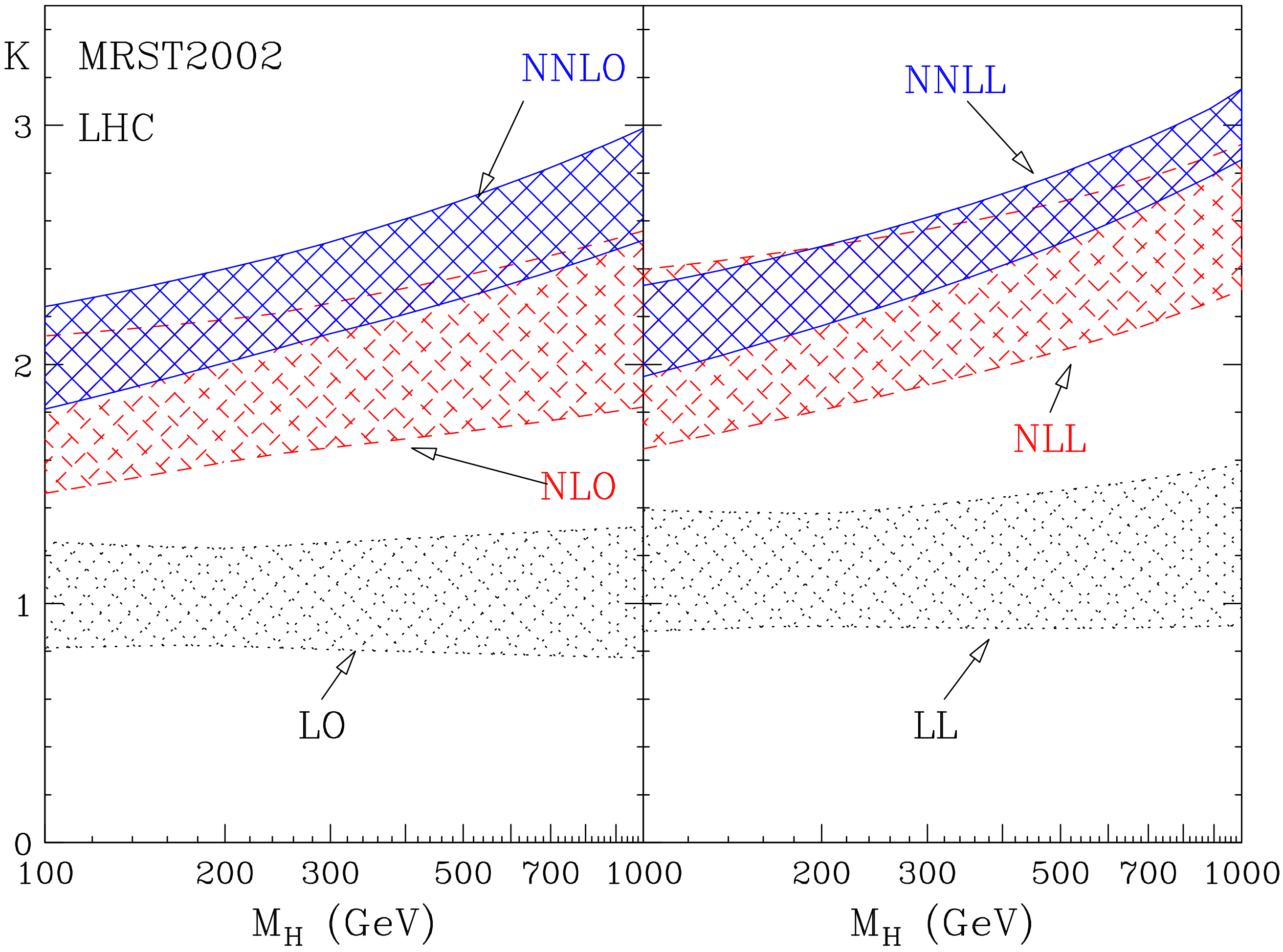
The function g_1 controls the LL contributions g_2 the NLL contributions g_3 the NNLL ones and so on

At NNLL three new coefficients appear:

- $D^{(2)}$, $C^{(2)}$ that can be obtained from the NNLO result
- $A^{(3)}$ which is known numerically

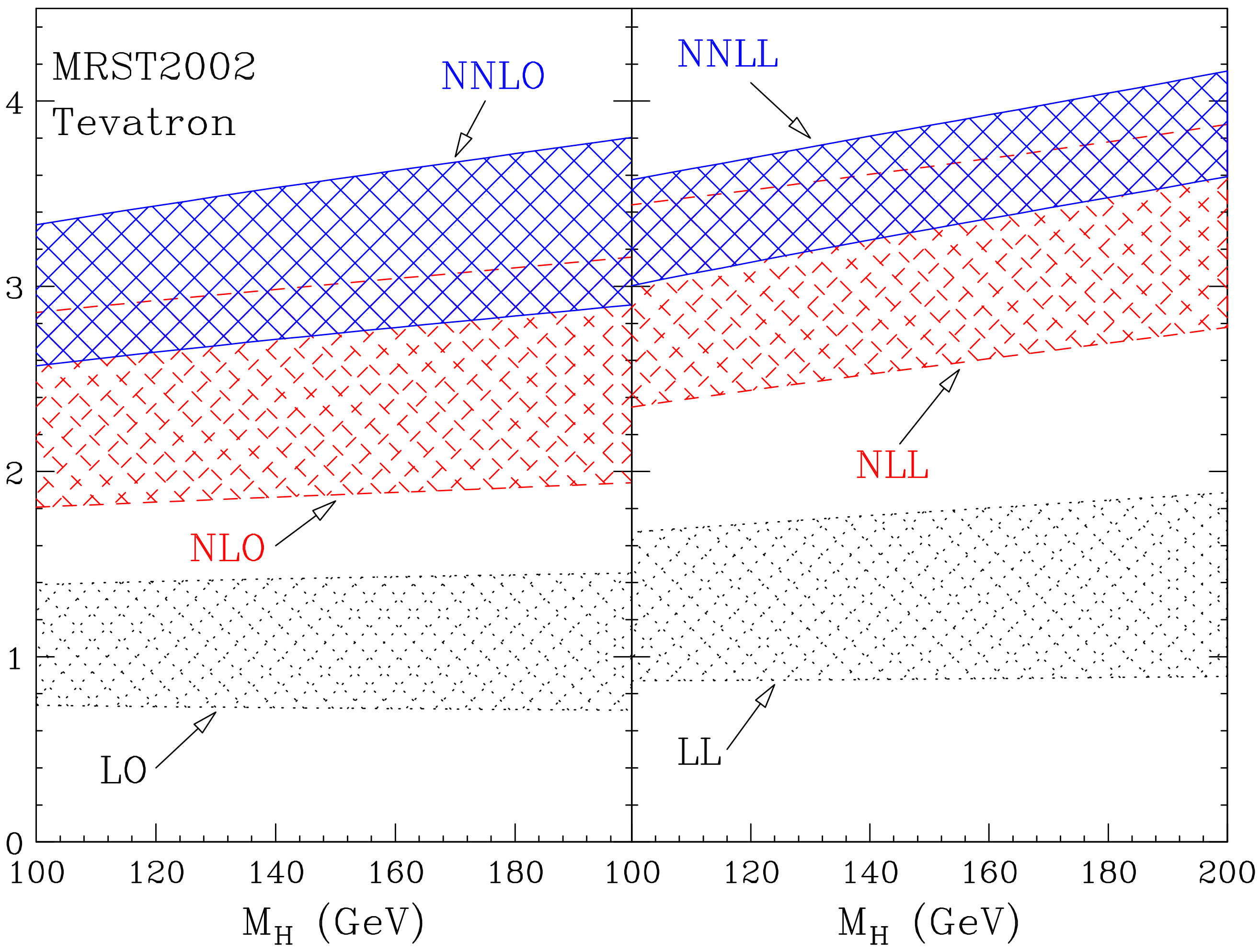
A.Vogt (2000)

Results at the LHC



- Resummed result matched to corresponding fixed order
- K-factors defined with respect to $\sigma^{LO}(\mu_F = \mu_R = M_H)$
  - with $\mu_{F(R)} = \chi_{F(R)} M_H$ and $1/2 \leq \chi_{F(R)} \leq 2$
  - but $1/2 \leq \chi_F/\chi_R \leq 2$
- For a light Higgs:
  - NNLL effect about +6% with respect to NNLO
  - Scale uncertanty at NNLL of about ±8%

Results at the Tevatron



- Effect of about $12-15\%$
- Bands defined as for LHC
- Reduction in scale uncertanty from about $\pm 13\%$ at NNLO to about $\pm 8\%$ at NNLL

What is the residual theoretical uncertainty on σ_H ?

- scale dependence

- large-m_{top} approximation
  - At NLO the approximation works well BECAUSE the cross section is dominated by soft radiation that is weakly sensible to the heavy quark loop
  - The dominance of soft contributions persists at NNLO $\Rightarrow$it is natural to expect the large-m_{top} approx. work well also at higher orders
  - Message from NLO:
    Use exact Born cross section
    (with m_{top}, m_b dependence) to normalize the result
  - Residual uncertainty from here $\lesssim$5%

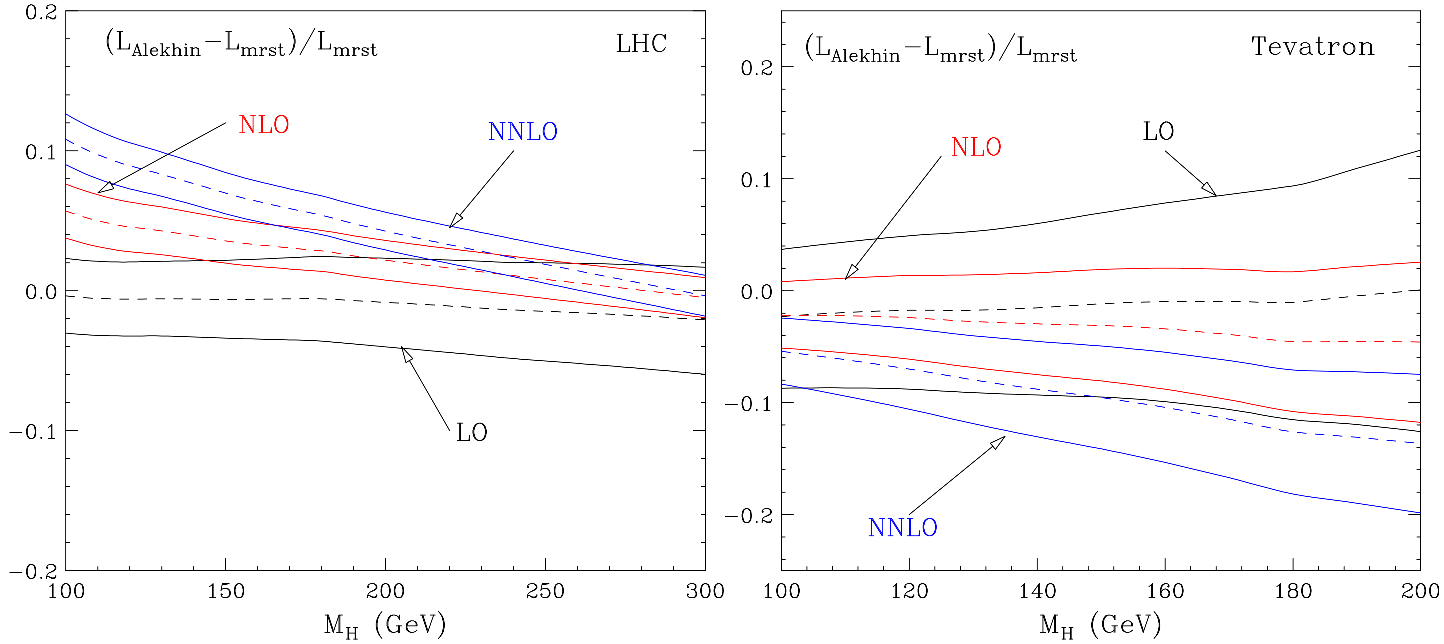
- Parton distributions
  - At NLO CTEQ6 and MRST2002 results agree reasonably well
  - At the moment only two approximated NNLO pdf sets: MRST, Alekhin

Comparing MRST and Alekhin results we find (relatively) large differences

- LHC: Alekhin results are larger than MRST : difference from 8% at $M_H = 100$ GeV to 2% at $M_H = 200$ GeV

- Tevatron: Alekhin results are smaller than MRST: difference from 7% at $M_H = 100$ GeV to 14% at $M_H = 200$ GeV

Errors only quoted by Alekhin and probably largely underestimated



The differences are due to the gg luminosities

Note: They increase with the order

The q_T spectrum of the Higgs

G. Bozzi, S. Catani, D. de Florian, MG (2003)

- Signal and background have different shape in q_T $\Rightarrow$knowledge of q_T distribution can help to improve statistical significance

- Higgs q_T spectrum is expected to be harder than $\gamma\gamma$ background

Studies of the Higgs q_T distribution have been performed at various levels of accuracy

I.Hinchliffe, S.F.Novaes (1988)
R.P.Kauffman (1992)
C.P Yuan (1992)
C.Balazs, C.P Yuan (2000)

Recently (almost) NNLL but still matching to LO

E.L.Berger, J.Qiu (2002)

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Our work:

- Include the most complete information which is available at present: NNLL resummation at small q_T and NLO perturbation theory at large q_T

- "Improve" the implementation formalism

Resummation

In the region $q_T^2 \ll Q^2$ large logarithmic corrections appear like $\alpha_S^n \log^{2n} Q^2/q_T^2$ that must be resummed to all orders

The q_T distribution can be decomposed as

$$\frac{d\sigma}{dq_T^2 dQ^2} = \frac{d\sigma^{(\text{res.})}}{dq_T^2 dQ^2} + \frac{d\sigma^{(\text{fin.})}}{dq_T^2 dQ^2}$$

- The first term contains all the logarithmically enhanced contributions and has to be resummed at all orders
- The finite part is not singular and can be evaluated at fixed order in PT

The resummation formalism has been developed in the eighties

G.Parisi, R.Petronzio (1979)
Y.Dokshitzer, D.Diakonov, S.I.Troian (1980)
G.Curci, M.Greco, Y.Srivastava (1979)
A.Bassetto, M.Ciafaloni, G.Marchesini (1980)
J.Kodaira, L.Trentadue (1982)
J.C.Collins, D.E.Soper, G.Sterman (1985)

As is customary in QCD resummations one has to work in a conjugate space in order to allow the kinematics of multiple gluon emission to factorize

In this case, to exactly implement transverse momentum conservation the resummation has to be performed in impact parameter b space

The resummation formula is usually written as

$$\frac{d\sigma^{(\text{res.})}}{dq_T^2 dQ^2} = \sum_{a,b,c} \int_0^1 dx_1 \int_0^1 dx_2 \int_0^\infty db \frac{b}{2} J_0(bq_T)\, \sigma_{c\bar{c}}^{(LO)}\, \delta(Q^2 - x_1 x_2 s)$$
$$\cdot \left(f_{a/h_1} \otimes C_{ca}\right) \left(x_1, \frac{b_0^2}{b^2}\right) \left(f_{b/h_2} \otimes C_{\bar{c}b}\right) \left(x_2, \frac{b_0^2}{b^2}\right) S_c(Q,b)$$

where $b_0 = 2e^{-\gamma}$ and $J_0(bq_T)$ have a kinematical origin

The large logarithmic corrections are exponentiated in the Sudakov form factor:

$$S_c(Q,b) = \exp\left\{ -\int_{b_0^2/b^2}^{Q^2} \frac{dq^2}{q^2} \left[A_c(\alpha_S(q^2))\, \ln \frac{Q^2}{q^2} + B_c(\alpha_S(q^2)) \right] \right\}$$

The coefficients A, B, C can be computed perturbatively

- The coefficients $A^{(1)}$, $A^{(2)}$ and $B^{(1)}$ are known both in the quark and the gluon channels

J.Kodaira, L.Trentadue (1982)
S.Catani, E. d'Emilio, L.Trentadue (1985)

- The coefficient $C^{(1)}$ is known for a variety of processes

C.Davies, W.J.Stirling (1984)
R.Kauffmann (1992)
C.Balazs, E.Berger, S.Mrenna, C.P.Yuan (1998)
D.de Florian, MG (2000)

- The general form of the NNLL coefficient $B^{(2)}$ has been computed recently

D. de Florian, MG (2000)

The standard “CSS” approach has several disadvantages:

a) The coefficients B and C are process dependent

D. de Florian, MG (2000)

$\Rightarrow$ For each process one is interested in (DY, Higgs, $\gamma\gamma$) new resummation coefficients have to be computed

b) The integral over the impact parameter b involves an extrapolation of the parton densities in the NP region

c) The resummation effects are large also at small b

- no control on the normalization
- problems in the matching with PT result
- unjustified higher order terms at large q_T with factorially growing coefficients (artifact of resummation)

S.Frixione, P.Nason, G.Ridolfi (1998)

To cure b) and c) resummation approaches directly in q_T space have been proposed

R.K.Ellis and S.Veseli (1998)
A.Kulesza and W.J.Stirling (1999)

$\Rightarrow$They can be at best approximate (no momentum conservation in transverse space)

Our formalism

A version of the b-space formalism has been proposed that overcomes all these problems

S. Catani, D. de Florian, MG (2000)

$$\frac{d\sigma}{dq_T^2}=\sum_{a,b}\int_0^1 dx_1\int_0^1 dx_2 f_{a/h_1}(x_1,\mu_F^2)f_{b/h_2}(x_2,\mu_F^2)\,\frac{d\hat\sigma_{ab}}{dq_T^2}$$

$$\frac{d\hat\sigma_{ab}}{dq_T^2}=\frac{d\hat\sigma_{ab}^{(\mathrm{res.})}}{dq_T^2}+\frac{d\hat\sigma_{ab}^{(\mathrm{fin.})}}{dq_T^2}$$

Parton distributions are evaluated at the fact. scale μ_F

$$\frac{d\hat\sigma_{ac}^{(\mathrm{res.})}}{dq_T^2}=\frac{1}{2}\int_0^\infty db\; b\; J_0(bq_T)\;\mathcal{W}_{ac}(b,M_H,\hat s;\alpha_S(\mu_R^2),\mu_R^2,\mu_F^2)$$

$$\begin{aligned}\mathcal{W}_N(b,M_H;\alpha_S(\mu_R^2),\mu_R^2,\mu_F^2)&=\mathcal{H}_N\left(\alpha_S(\mu_R^2);M_H^2/\mu_R^2,M_H^2/\mu_F^2\right)\\&\times\exp\{\mathcal{G}_N(\alpha_S(\mu_R^2),bM_H;M_H^2/\mu_R^2,M_H^2/\mu_F^2)\}\end{aligned}$$

where flavour indices should be understood

Then we take $\mu_F\sim\mu_R\sim M_H$ and organize the large logs as

$$\begin{aligned}\mathcal{G}_N(\alpha_S,bM_H;M_H^2/\mu_R^2,M_H^2/\mu_F^2)&=L\,g^{(1)}(\alpha_S L)+g_N^{(2)}(\alpha_S L;M_H^2/\mu_R^2)\\&+\alpha_S\,g_N^{(3)}(\alpha_S L;M_H^2/\mu_R^2,M_H^2/\mu_F^2)+\dots\end{aligned}$$

where $L=\ln M_H^2b^2/b_0^2$ and $\alpha_S=\alpha_S(\mu_R)$

The form factor takes thus the same structure as in e^+e^- event shape variables or in threshold resummation in hadron collisions

⇒A study of scale dependence can be performed as is normally done in fixed order calculation (no need of introducing additional coefficients)

The functions $g^{(n)}(\lambda)$ are defined such that $g^{(n)}(0) = 0$ and:

- $g^{(1)}$ depends on $A^{(1)}$ (LL)
- $g_N^{(2)}$ depends on $A^{(1)}$, $B^{(1)}$ and $A^{(2)}$ (NLL)
- $g_N^{(3)}$ also depends on $C^{(1)}$, $B^{(2)}$ and $A^{(3)}$ (NNLL)

The functions $g_N^{(2)}$ and $g_N^{(3)}$ respectively receive additional contribution from the LO and NLO anomalous dimensions

Notice: Since the large log is $L = \ln\left(\frac{M_H^2 b^2}{b_0^2}\right)$ the form factor is divergent as $b \to 0$

$\Rightarrow$we perform the replacement

$$L \longrightarrow \widetilde{L} \equiv \ln\left(1 + \frac{M_H^2 b^2}{b_0^2}\right)$$

- This replacement is legitimate in the large-b region where $\widetilde{L} \sim L$
- At small b the resummation effects vanish ($\widetilde{L} \to 0$)
- The replacement is inspired by the procedure followed in e^+e^- event shapes

S.Catani, L.Trentadue, G.Turnock, B.R.Webber (1993)

Landau pole avoided by diverting the b-space integration to the complex plane

E.Laenen, G.Sterman, W.Vogelsang (2000)

Numerical results

I present NLL results matched to LO (NLL+LO) and NNLL results matched to NLO (NNLL+NLO)

We use MRST2001 pdf

- NLL+LO: LO pdf + 1-loop α_S
- NNLL+NLO: NLO pdf + 2-loop α_S

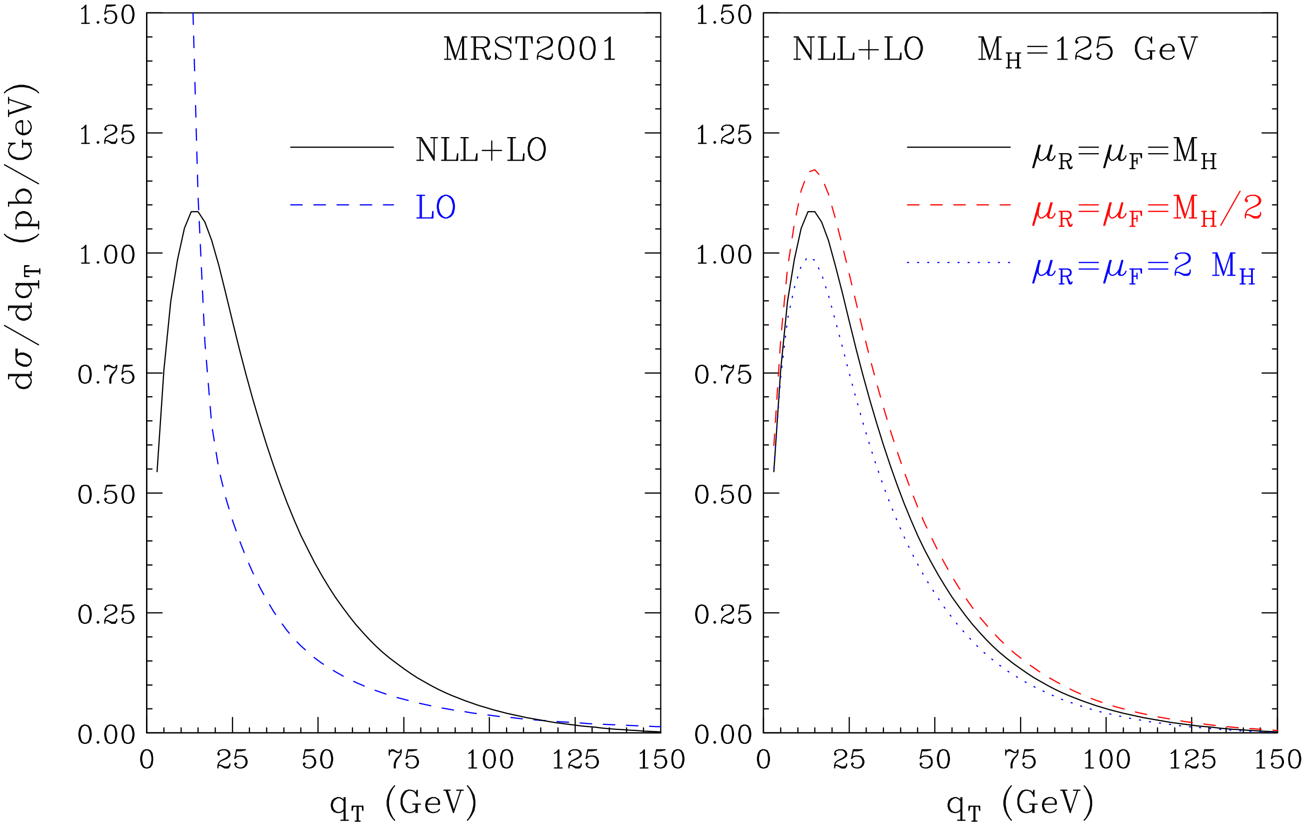
At NNLL+NLO $A^{(3)}$ and $\mathcal{H}^{(2)}$ are not known

- For $A^{(3)}$ we use the numerical estimate available for threshold resummation

A. Vogt (2000)

- The effect of $\mathcal{H}^{(2)}$ is included in approximated form using the known result for the total NNLO cross section (computed with NLO pdf and 2-loop α_S)

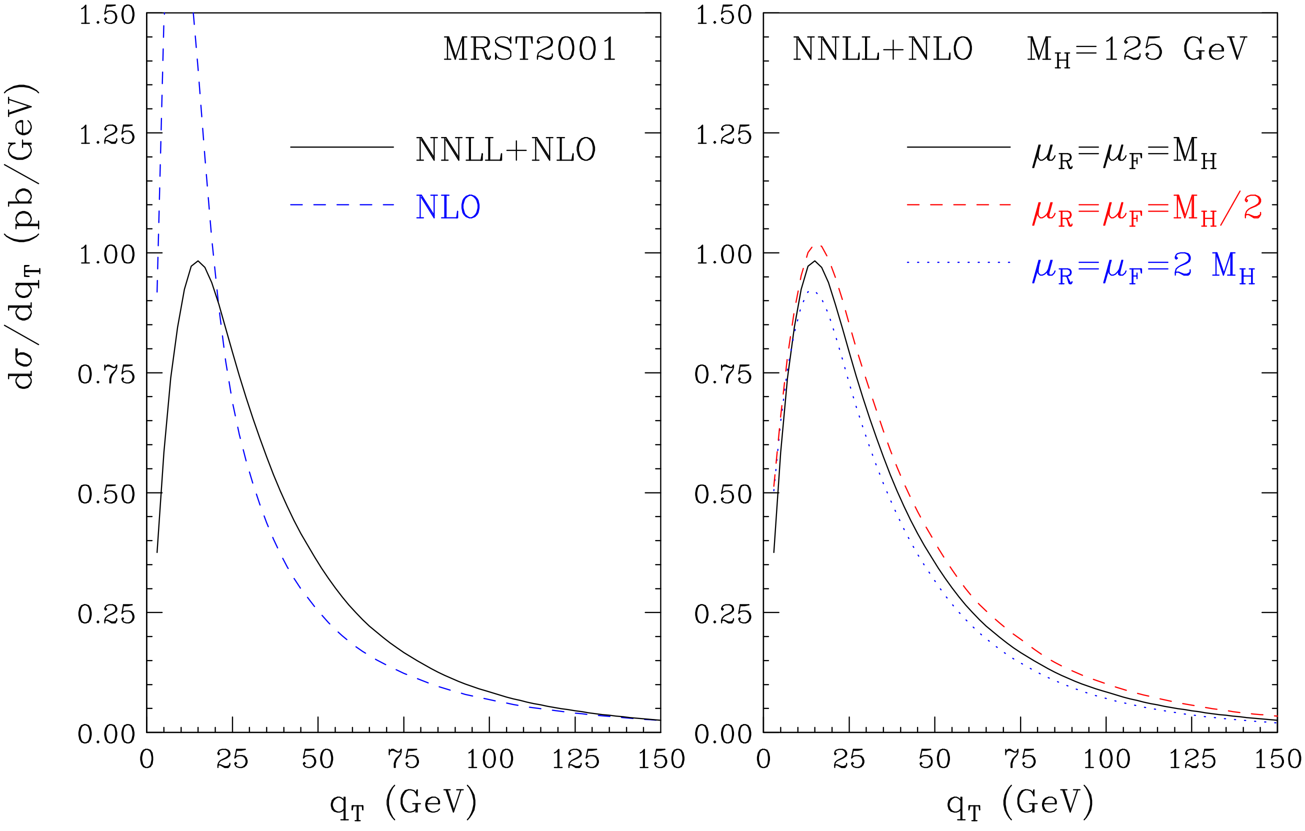
CTEQ6 pdf set gives very similar results



The LO result diverges to $+\infty$ as $q_T \to 0$

The effect of resummation is relevant below $q_T \sim 100$ GeV

Total cross section in good agreement with NLO result computed with LO pdf and 1-loop α_S



The NLO result diverges to $-\infty$ as $q_T \to 0$ (unphysical peak)

The q_T distribution is slightly harder than at NLL+LO ($\langle q_T \rangle \sim 39$ GeV)

The effect of the coefficient $A^{(3)}$ is negligible whereas the coefficient $\mathcal{H}^{(2)}$ increases the cross section by about 20%

The scale dependence is reduced with respect to NLL+LO: It varies from $\pm 10\%$ at the peak to $\pm 20\%$ at $q_T \sim 100$ GeV

Summary

We have evaluated the contribution of multiple soft-gluon emissions to the the total cross section σ_H:

- Effect moderate at LHC: for a light Higgs $+6\%$ with respect to NNLO
- A bit larger at the Tevatron: $+12-15\%$ with respect to NNLO

$\Rightarrow$ Perturbative result under better control now but...
still problems with NNLO pdf !

We have computed the q_T spectrum of the SM Higgs boson at the LHC

- We have implemented the most complete information available at present:
  all-order resummation of logarithmically enhanced contributions at small q_T at NNLL level combined with NLO perturbation theory at large q_T
- Distinctive feature of our approach are:
  - It allows a consistent study of theoretical uncertainties
  - It avoids the introduction of unjustified higher-order corrections in the intermediate-q_T region by using unitarity constraints $\Rightarrow$Normalization OK
- The results appear to be stable
- The method can be used also with other processes (DY,$\gamma\gamma$...)