

Isolated γ production in the POWHEG + MiNLO framework: the $W\gamma$ example

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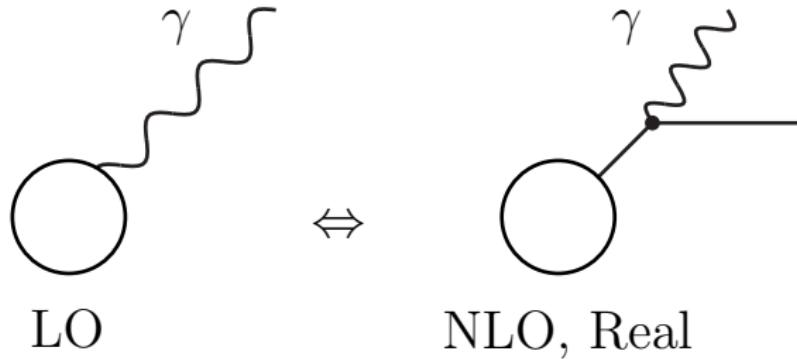
INFN Sezione di Pavia

June 6th 2015, Les Houches

in collaboration with L. Barzè, G. Montagna, P. Nason,
O. Nicrosini, F. Piccinini and V. Prospieri

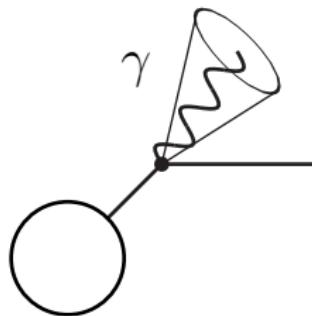
based on JHEP **1412**, 039 (2014)

NLO QCD corrections with isolated γ



- QED singularities from the integration of real QCD corrections
- no virtual counterpart, ...

Theory wayout (1): Frixione isolation¹



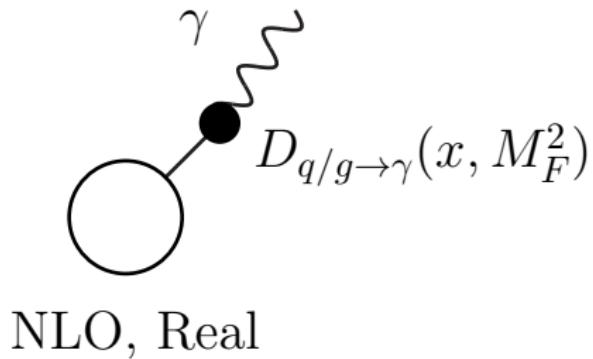
NLO, Real

$$\forall_{R < R_0} \sum_{R_{j,\gamma} < R} E_{T,j} < \epsilon_h p_T^\gamma \left(\frac{1 - \cos R}{1 - \cos R_0} \right)$$

not directly comparable
with data

¹S. Frixione, Phys. Lett. B **429** (1998) 369

Theory wayout (2): fragmentation functions²



- extracted from data
- **collinear** approximation

²S. Catani et al., JHEP 0205 (2002) 028

Theory wayout (3): remark

both fragmentation functions and frixione isolation:

- useful and well defined for fixed order calculations
- impose a separation between final state partons and photons

$W\gamma$ production in POWHEG: target

- predictions for $pp \rightarrow W(l\nu)\gamma$
at NLO+PS QCD accuracy
- exclusive generation of radiated particles
(q/g or γ) in the POWHEG framework

Theoretical predictions: state of the art (1)

- NLO QCD corrections:
J. Ohnemus Phys. Rev. **D47** (1993),
U. Baur et al., Phys. Rev. **D48** (1993) and Phys. Rev. **D57** (1998)
D. De Florian and A. Signer Eur. Phys. J. **C16** (2000)
- NLO EW corrections (leading pole):
E. Accomando et al., Eur. Phys. J. **C47** (2006)
- NLO EW and QCD corrections:
A. Denner et al., JHEP **1504**, 018 (2015)
- NNLO QCD corrections:
M. Grazzini arXiv:1504.01330
M. Grazzini et al., Phys.Lett. **B731** (2014),
M. Grazzini arXiv:1407.1618

Theoretical predictions: state of the art (2)

isolated γ production in MC event generators

- LO matrix element merged with QCD+QED PS
(isolated γ and $\gamma\gamma$ production):
S. Hoeche et al., Phys. Rev. **D81** (2010)
- $\gamma\gamma$ production with the POWHEG method:
L. D'Errico and P. Richardson JHEP **1202** (2012)
- $t\bar{t}\gamma$ and $t\bar{t}\gamma\gamma$ production in the POWHEG BOX framework:
A. Kardos and T. Trocsany arXiv:1406.2324, arXiv:1408.0278

POWHEG³ in a nutshell: QCD processes (1)

$$\begin{aligned} d\sigma = & \sum_{f_b} \bar{B}^{f_b}(\Phi_n) d\Phi_n \left\{ \Delta^{f_b}(\Phi_n, p_T^{\min}) \right. \\ & + \left. \sum_{\alpha_r \in \{\alpha_r | f_b\}} \frac{\left[d\Phi_{\text{rad}} \theta(k_T - p_T^{\min}) \Delta^{f_b}(\Phi_n, k_T) R(\Phi_{n+1}) \right]_{\alpha_r}^{\bar{\Phi}_n^{\alpha_r} = \Phi_n}}{B^{f_b}(\Phi_n)} \right\} \end{aligned}$$

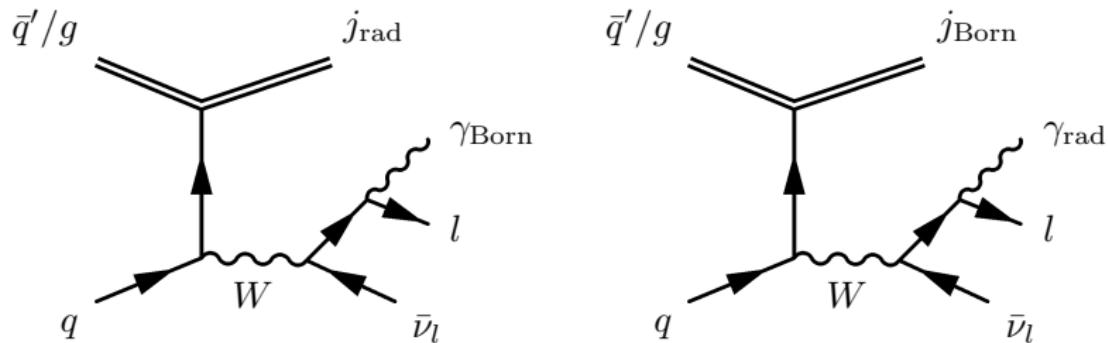
³P. Nason, JHEP **0411** (2004), S. Frixione et al., JHEP **0711** (2007),
S. Alioli et al., JHEP **1006** (2010)

POWHEG in a nutshell: QCD processes (2)

$$\begin{aligned}
 \bar{B}^{f_b}(\Phi_n) &= [B(\Phi_n) + V(\Phi_n)]_{f_b} \\
 &+ \sum_{\alpha_r \in \{\alpha_r | f_b\}} \int \left[d\Phi_{\text{rad}} \{R(\Phi_{n+1}) - C(\Phi_{n+1})\} \right]_{\alpha_r}^{\bar{\Phi}_n^{\alpha_r} = \Phi_n} \\
 &+ \sum_{\alpha_{\oplus} \in \{\alpha_{\oplus} | f_b\}} \int \frac{dz}{z} G_{\oplus}^{\alpha_{\oplus}}(\Phi_{n,\oplus}) + \sum_{\alpha_{\ominus} \in \{\alpha_{\ominus} | f_b\}} \int \frac{dz}{z} G_{\ominus}^{\alpha_{\ominus}}(\Phi_{n,\ominus})
 \end{aligned}$$

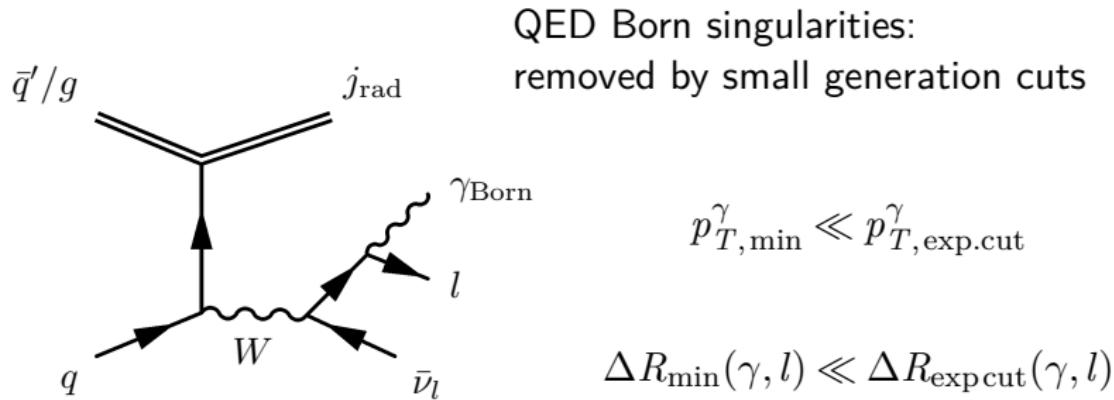
$$\begin{aligned}
 \Delta^{f_b}(\Phi_n, p_T) &= \\
 \exp \left\{ - \sum_{\alpha_r \in \{\alpha_r | f_b\}} \int \frac{\left[d\Phi_{\text{rad}} R(\Phi_{n+1}) \theta(k_T(\Phi_{n+1}) - p_T) \right]_{\alpha_r}^{\bar{\Phi}_n^{\alpha_r} = \Phi_n}}{B^{f_b}(\Phi_n)} \right\}
 \end{aligned}$$

$W\gamma$ production in POWHEG: starting point



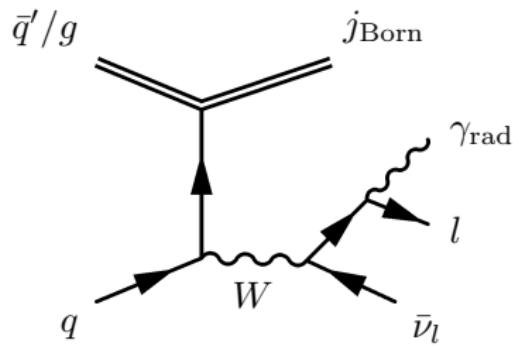
- QCD (QED) singular regions are separated (FKS method) and mapped on the corresponding underlying Born configurations $W\gamma$ (Wj)
- hardest emission (q/g or γ) generated from the modified POWHEG Sudakov form factor

$W\gamma$ production in POWHEG: $W\gamma$ Underlying Born



Born $W\gamma \times$ QCD Rad.

$W\gamma$ production in POWHEG: Wj Underlying Born



QCD Born singularities:

- no exp. cuts on $j_{\text{Born}} \Rightarrow$ no safe generation cuts on j_{Born}
- MiNLO⁴ provides the Sudakov suppression factor that cancels the Born QCD singularities

Born $Wj \times$ QED Rad.

⁴K. Hamilton et al., JHEP **1210** (2012),
K. Hamilton et al., JHEP **1305** (2013)

$W\gamma$ production in POWHEG: remark

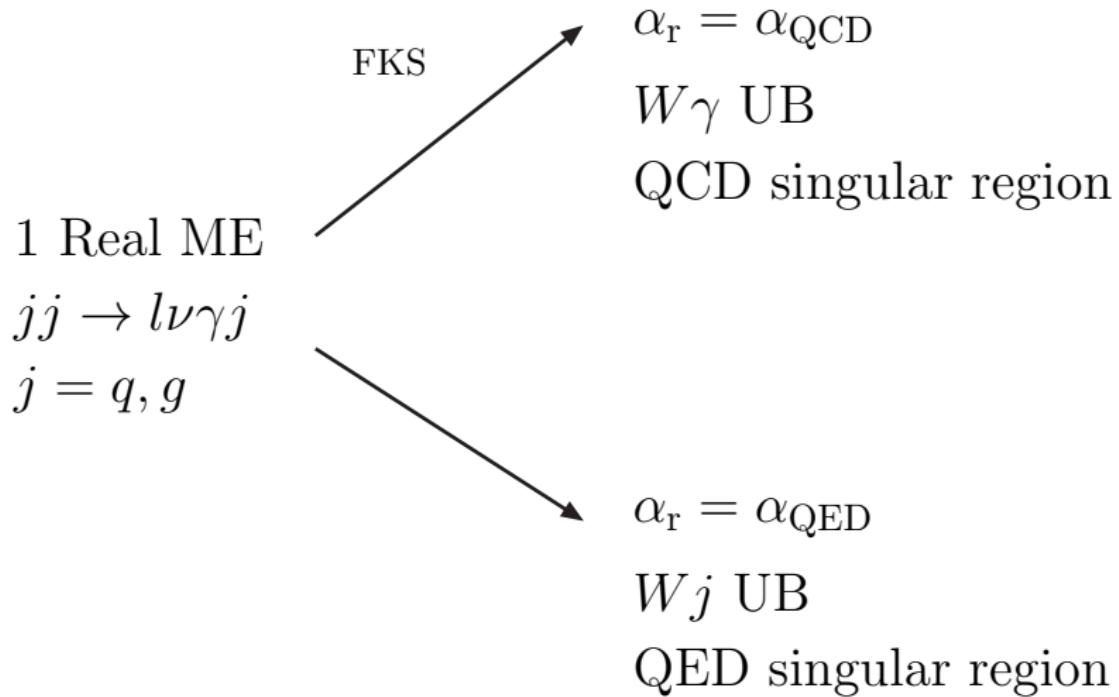
- MiNLO is mandatory for Wj underlying Born
- for $W\gamma$ underlying Born MiNLO is used in order to improve the α_S scale choice in the QCD corrections

$W\gamma$ production in POWHEG: implementation

- NC *with No Competition*
- C-LO *with Competition, LO normalization*
- C-NLO *with Competition, NLO normalization*

Competition
between QCD and QED radiation off quarks

$W\gamma$ production in POWHEG: NC (1)



$W\gamma$ production in POWHEG: NC (2)

for $W\gamma$ underlying Born⁵

- normalization:

$$\overline{B}_{W\gamma} = B_{W\gamma} + V_{W\gamma} + \sum_{\alpha_{\text{QCD}}} \int [d\Phi_{\text{Rad}}(R_{W\gamma;j} - C)]^{\alpha_{\text{QCD}}}$$

- radiation dynamics:

$$\Delta_{W\gamma}^{\alpha_{\text{QCD}}} = \exp\left\{-\int [d\Phi_{\text{Rad}} \frac{R_{W\gamma;j}}{B_{W\gamma}}]^{\alpha_{\text{QCD}}}\right\}$$

- scalup choice (POWHEG standard):

$$p_T^{\text{Rel.}}(j_{\text{rad.}})$$

⁵simplified notation: MiNLO factors understood

$W\gamma$ production in POWHEG: NC (3)

for Wj underlying Born

- normalization:

$$\overline{B}_{Wj} = B_{Wj} + \sum_{\alpha_{\text{QED}}} \int [d\Phi_{\text{Rad}}(R_{Wj;\gamma} - C)]^{\alpha_{\text{QED}}}$$

- radiation dynamics:

$$\Delta_{Wj}^{\alpha_{\text{QED}}} = \exp\left\{-\int [d\Phi_{\text{Rad}} \frac{R_{Wj;\gamma}}{B_{Wj}}]^{\alpha_{\text{QED}}}\right\}$$

- scalup choice:

- $p_T(j_{\text{Born}})$ for QCD shower
- $p_T^{\text{Rel.}}(\gamma_{\text{rad.}})$ for QED shower

$W\gamma$ production in POWHEG: C-LO

$W\gamma$ underlying Born: same as NC

for Wj underlying Born

- normalization: same as for NC
- radiation dynamics:

$$\Delta_{Wj} = \exp \left\{ - \sum_{\alpha_{\text{QED}}} \int [d\Phi_{\text{Rad}} \frac{R_{Wj;\gamma}}{B_{Wj}}]^{\alpha_{\text{QED}}} \right. \\ \left. - \sum_{\alpha_{\text{QCD}}} \int [d\Phi_{\text{Rad}} \frac{R_{Wj;j}}{B_{Wj}}]^{\alpha_{\text{QCD}}} \right\}$$

- scalup choice (POWHEG standard): $p_T^{\text{Rel}}(\text{rad})$

$W\gamma$ production in POWHEG: C-NLO

- same dynamics of C-LO
- NLO Wj normalization

$$\begin{aligned}\overline{B}_{Wj} = & B_{Wj} + V_{Wj}^{\text{QCD}} + \sum_{\alpha_{\text{QED}}} \int [d\Phi_{\text{Rad}}(R_{Wj;\gamma} - C)]^{\alpha_{\text{QED}}} \\ & + \sum_{\alpha_{\text{QCD}}} \int [d\Phi_{\text{Rad}}(R_{Wj;j} - C)]^{\alpha_{\text{QCD}}}\end{aligned}$$

Comparison with ATLAS data⁶: event selection

$$p_T^\gamma > 15 \text{ GeV}, |\eta_\gamma| < 2.37, \Delta R_{\ell\gamma} > 0.7,$$

$$R_0 = 0.4, \epsilon_h = 0.5, \sum_{R_{j,\gamma} < R_0} E_{T,j} < \epsilon_h p_T^\gamma$$

$$p_T^\ell > 25 \text{ GeV}, |\eta_\ell| < 2.47, p_T^{\text{miss.}} > 35 \text{ GeV}$$

$$E_T^{\text{jet}} > 30 \text{ GeV}, |\eta_{\text{jet}}| < 4.4, \Delta R(e/\mu/\gamma, \text{jet}) > 0.3$$

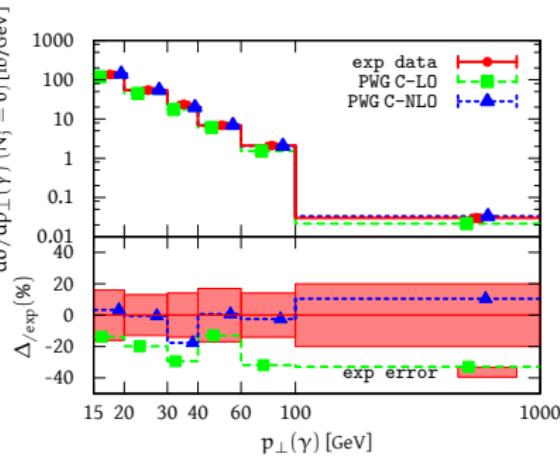
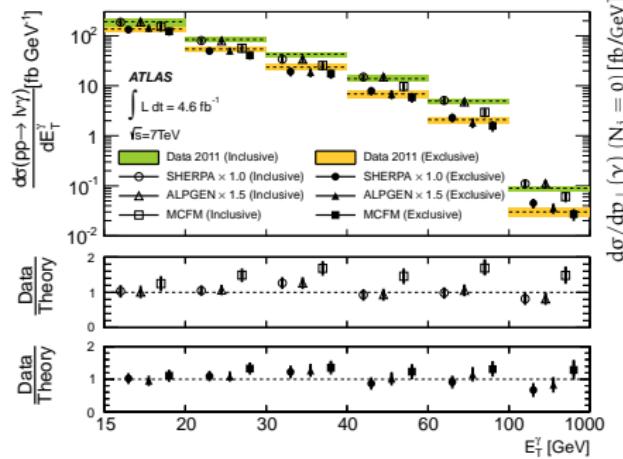
⁶Phys. Rev. **D87** (2013)

Comparison with ATLAS data: cross sections

[Pb]	$N_{\text{jet}} = 0$	$N_{\text{jet}} \geq 0$
Exp.	$1.77^{+0.04}_{-0.08} \text{stat} \pm 0.24 \text{syst}$	$2.74^{+0.05}_{-0.14} \text{stat} \pm 0.32 \text{syst}$
MCFM	1.39 ± 0.17	1.96 ± 0.17
C-LO	$1.42^{+0.15}_{-0.15}$	$2.25^{+0.24}_{-0.24}$
C-NLO	$1.69^{+0.11}_{-0.22}$	$2.95^{+0.20}_{-0.38}$

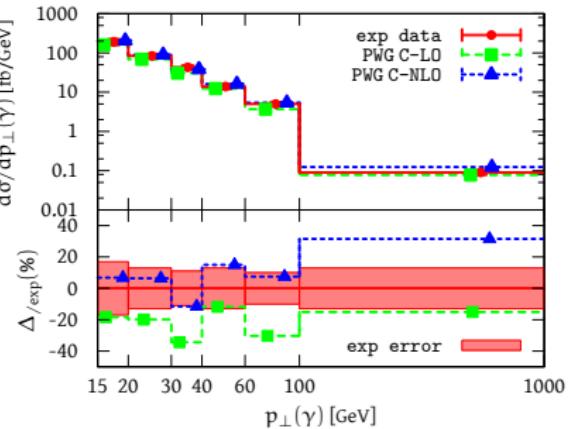
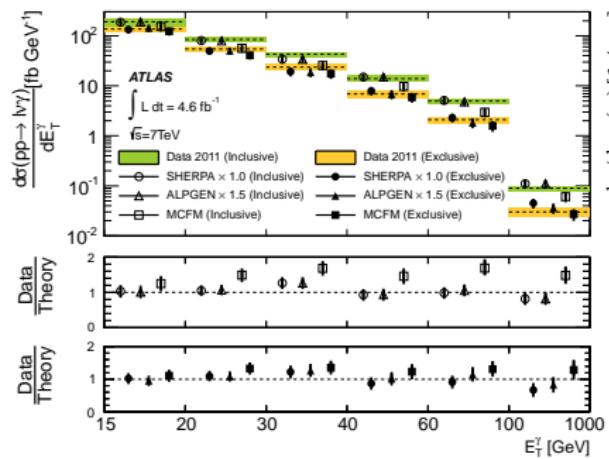
Comparison with ATLAS data: distributions (1)

p_T^γ , $N_{\text{jet}} = 0$

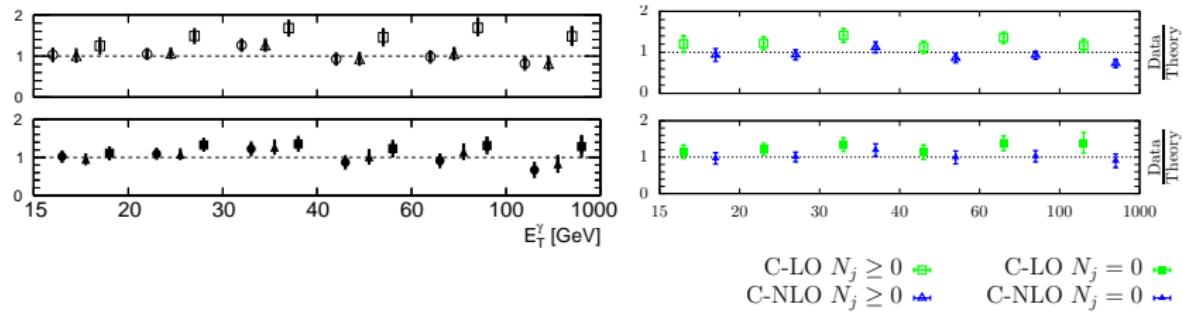


Comparison with ATLAS data: distributions (2)

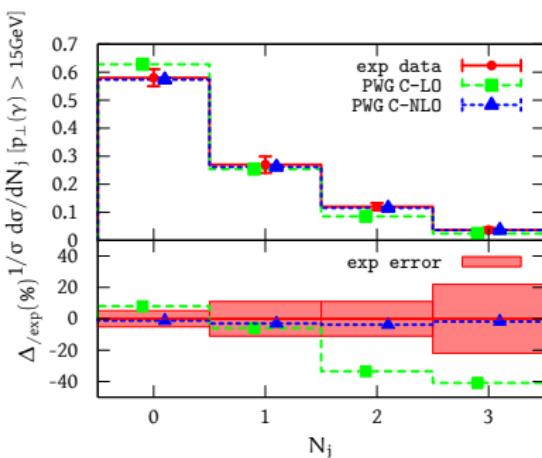
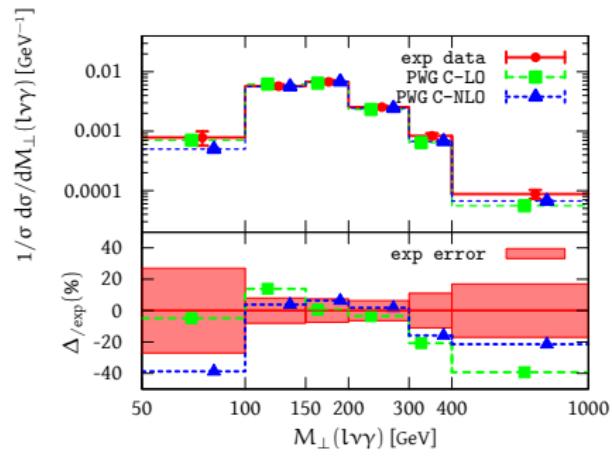
$p_T^\gamma, N_{\text{jet}} \geq 0$



Comparison with ATLAS data: distributions (3)

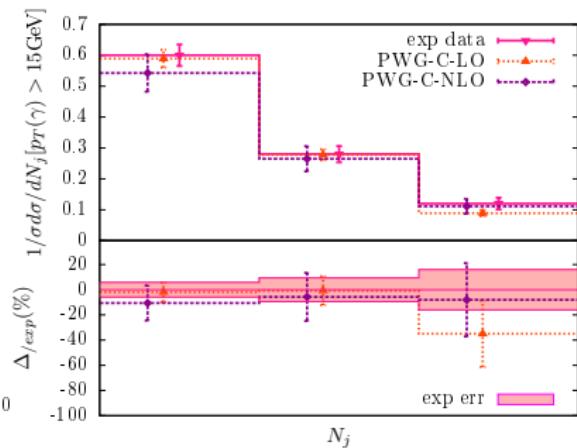
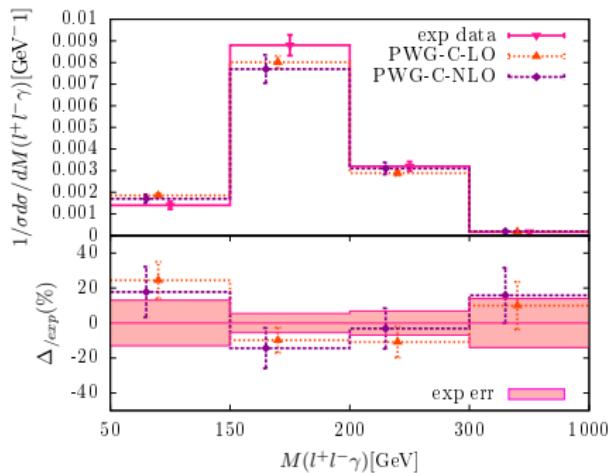


Comparison with ATLAS data: distributions (4)



Conclusions

- $W\gamma$ production process implemented in the POWHEG-BOX-V2 generator
- general treatment of final states isolated photons: $Z\gamma$ (preliminary)



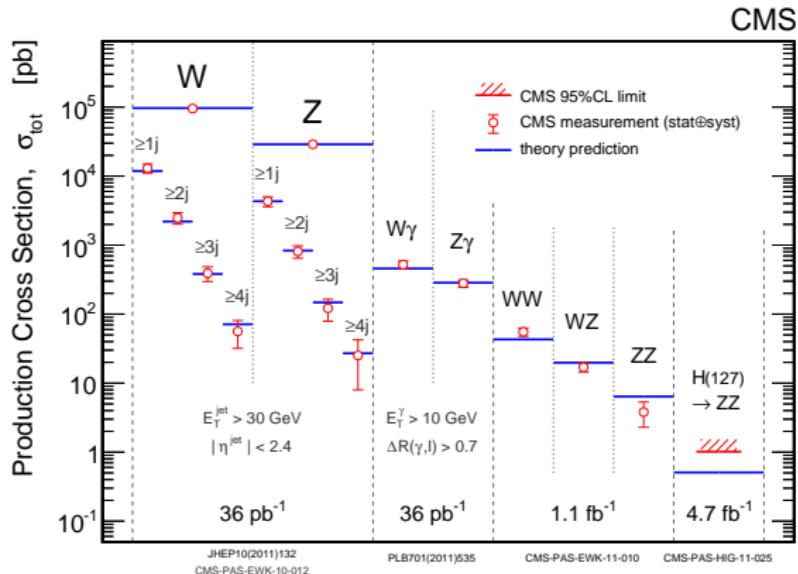
by Valeria Prosperi

Conclusions

- $W\gamma$ production process implemented in the POWHEG-BOX-V2 generator
 - NLO QCD normalization
 - exclusive generation of the final state particles
- general treatment of isolated photons in the final states: $Z\gamma$
- *no need to impose generation cuts*
- improvement in the data/theory comparison

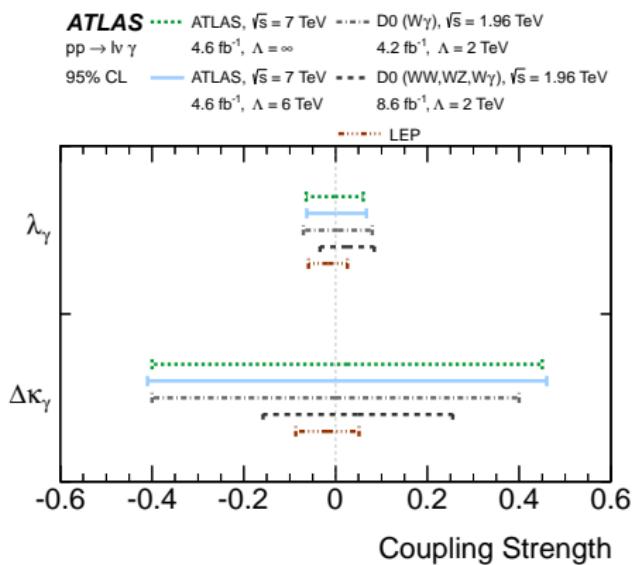
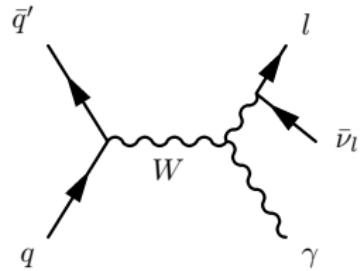
Backup Slides

Motivations (1): luminosity⁷



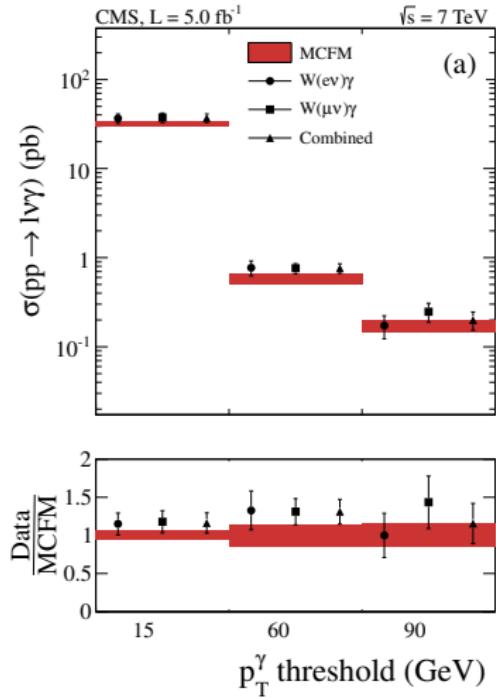
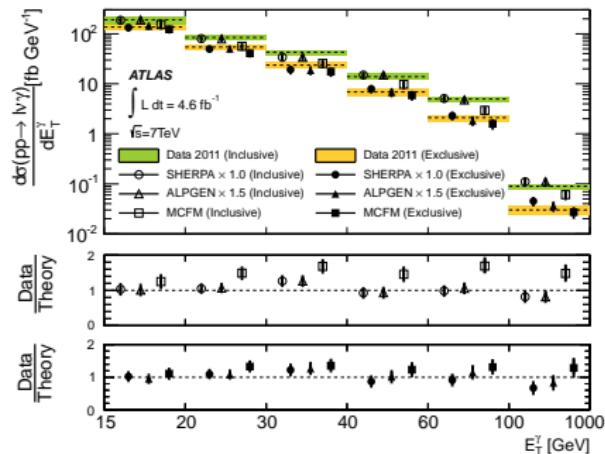
⁷from CMS public results (a bit old)

Motivations (2): ATGCs⁸



⁸from ATLAS public results

Motivations (3): Data VS Theory (tensions)⁹



⁹G. Aad et al. Phys. Rev. **D87** (2013),
S. Chatrchyan et al. Phys. Rev. **D89** (2013)

Comparison with MCFM, event selections¹⁰

- Basic Photon cuts:

$$p_T^\gamma > 15 \text{ GeV}, |\eta_\gamma| < 2.37, \Delta R_{\ell\gamma} > 0.7,$$
$$R_0 = 0.4, \epsilon_h = 0.5$$

- $M_T(l\nu\gamma)$ cuts:

Basic Photon + $M_T > 90$ GeV

- Lepton cuts:

Basic Photon + $p_T^\ell > 25$ GeV, $|\eta_\ell| < 2.47$, $p_T^\nu > 35$ GeV

¹⁰as in JHEP **1107** (2011) 018

Comparison with MCFM: NLO results

Cuts	MCFM	POWHEG NLO
Basic Photon	13.12(4)	13.15(1)
M_T cut	2.770(1)	2.774(3)
Lepton cuts	1.126(1)	1.123(4)

7 TeV

Cuts	MCFM	POWHEG NLO
Basic Photon	23.90(8)	24.04(3)
M_T cut	6.230(2)	6.250(9)
Lepton cuts	2.342(2)	2.340(6)

14 TeV

Frixione isolation procedure

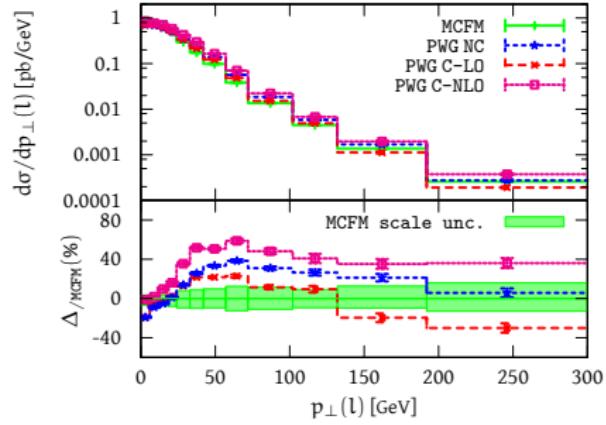
Comparison with MCFM: full simulation (1)

7 TeV				
Cuts	MCFM	POWHEG-NC	C-LO	C-NLO
Basic	$12.92(3)^{+4\%}_{-6\%}$	$13.11(3)^{+8\%}_{-10\%}$	$12.95(3)^{+8\%}_{-11\%}$	$15.08(7)^{+2\%}_{-9\%}$
M_T	$2.625(1)^{+6\%}_{-6\%}$	$3.65(2)^{+10\%}_{-11\%}$	$3.20(2)^{+11\%}_{-11\%}$	$4.23(3)^{+10\%}_{-10\%}$
Lept.	$1.077(1)^{+6\%}_{-6\%}$	$1.44(1)^{+8\%}_{-10\%}$	$1.31(1)^{+11\%}_{-11\%}$	$1.75(2)^{+7\%}_{-13\%}$

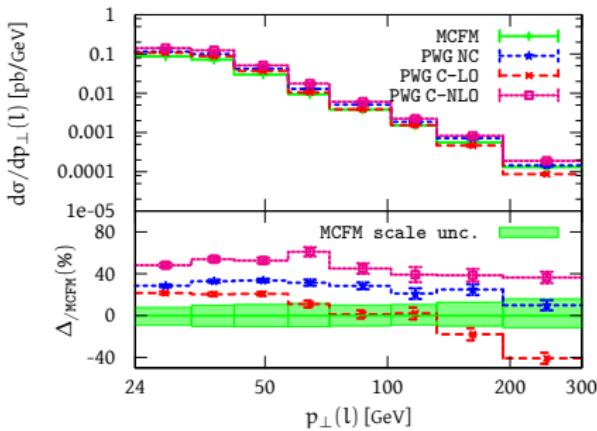
- MCFM: fragmentation funtions
- POWHEG: experimental-like isolation condition

$$\sum_{R_{j,\gamma} < R_0} E_{T,j} < \epsilon_h p_T^\gamma$$

Comparison with MCFM: full simulation (2)

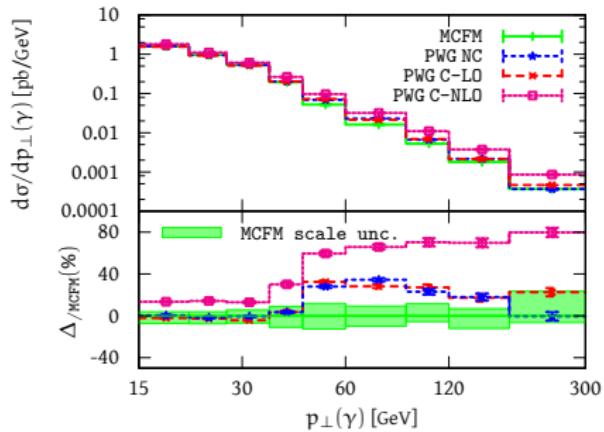


Basic cuts

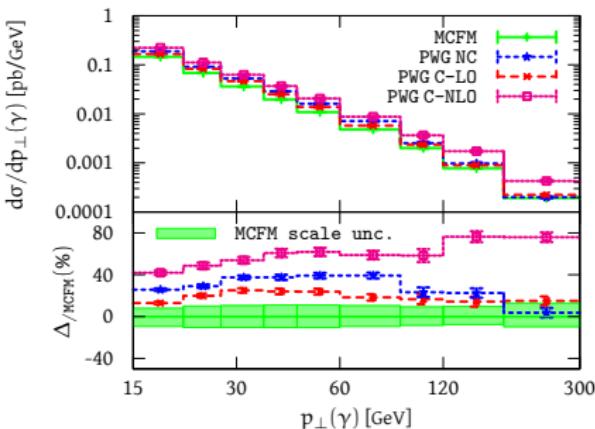


Lepton cuts

Comparison with MCFM: full simulation (3)



Basic cuts



Lepton cuts