

# **$q_T$ subtraction for heavy quark production**

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# $q_T$ subtraction - colourless FS

- Originally developed for the production of colourless final state  $F$  in hadron collisions
- Slicing method, slicing parameter:  $q_T$  (transverse momentum of system  $F$ )

$$d\sigma_{\text{NNLO}}^F \Big|_{q_T \neq 0} = d\sigma_{\text{NLO}}^{F+\text{jet}} \quad \begin{array}{l} \xrightarrow{\text{RR and RV contributions from } F@ \text{NNLO}} \\ \xrightarrow{\text{Computable with any NLO subtraction method, IR finite}} \end{array}$$

- Only  $q_T \rightarrow 0$  infrared divergencies remain
- But small  $q_T$  behavior known from  $q_T$  resummation

$$d\sigma_{\text{NNLO}}^F = \mathcal{H}_{\text{NNLO}}^F \otimes d\sigma_{\text{LO}}^F + \underbrace{\left[ d\sigma_{\text{NLO}}^{F+\text{jet}} - d\sigma_{\text{NNLO}}^{\text{CT}} \right]}_{\text{Finite for } q_T \rightarrow 0}$$

$\downarrow$   
 Hard-collinear contributions,  
 only at  $q_T=0$

Restores correct  
normalization, includes the  
2-loop corrections

# $q_T$ subtraction - colourless FS

Finite for  $q_T \rightarrow 0$

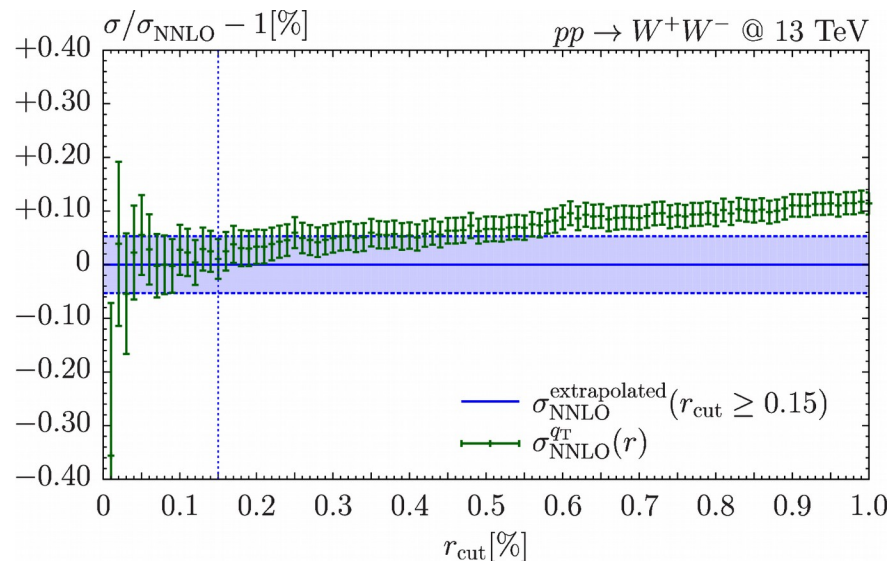
$$d\sigma_{\text{NNLO}}^F = \mathcal{H}_{\text{NNLO}}^F \otimes d\sigma_{\text{LO}}^F + \left[ d\sigma_{\text{NLO}}^{F+\text{jet}} - d\sigma_{\text{NNLO}}^{\text{CT}} \right]$$

We need to introduce a cutoff

Individually divergent when  $q_T \rightarrow 0$

$$r = q_T/Q$$

- Introduce an  $r_{\text{cut}}$
- Check that results are independent of  $r_{\text{cut}}$
- Extrapolate  $r_{\text{cut}} \rightarrow 0$  result



- $q_T$  subtraction successfully applied to a large number of processes (V,  $VV'$ , H, HH)
- Public implementation for fully differential calculations: MATRIX  
[Grazzini, Kallweit, Wiesemann, '17]

# $q_T$ subtraction - colourless FS

$q_T$  subtraction “solved” at NNLO for colourless final-states

( First N3LO application  
in 1807.11501 for H prod )

$$d\sigma_{\text{NNLO}}^F = \mathcal{H}_{\text{NNLO}}^F \otimes d\sigma_{\text{LO}}^F + \left[ d\sigma_{\text{NLO}}^{F+\text{jet}} - d\sigma_{\text{NNLO}}^{\text{CT}} \right]$$

Obtained from the  
hard-collinear function

$$[H C_1 C_2]$$

Universal relation between  
hard-virtual function and  
loop-corrections known

$$H \sim \langle \tilde{\mathcal{M}} | \tilde{\mathcal{M}} \rangle \quad \text{with} \quad |\tilde{\mathcal{M}}\rangle = \left( 1 - \tilde{I} \right) |\mathcal{M}\rangle$$

All-orders loop amplitude  
(process dependent)

$$d\sigma^{\text{CT}} \sim \Sigma^F(q_T/M) \otimes d\sigma_{\text{LO}}^F$$

CT coefficients known

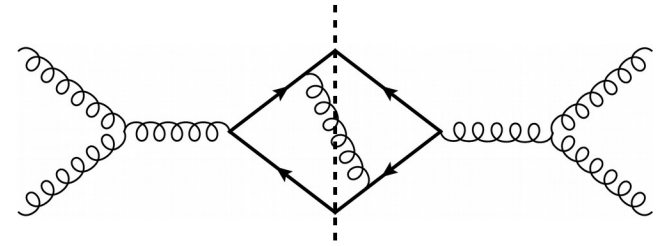
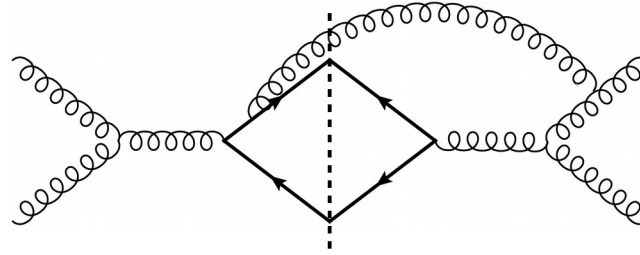
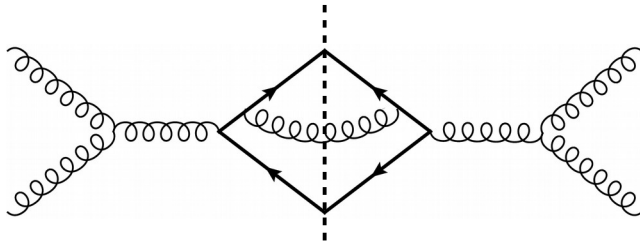
Subtraction operator that correctly  
accounts for the stuff we put in the CT  
(universal)

Computing  $H \longleftrightarrow$  Computing  $\tilde{I}$  (+ virtuals)

Can be explicitly obtained by computing the differential cross  
section  $d\sigma/d^2q_T$  at small  $q_T$  at the corresponding perturbative order

# $q_T$ subtraction - $Q\bar{Q}$ case

Additional final-state radiation  $\longrightarrow$  Massive emitters: only new **soft** divergencies



# $q_T$ subtraction - $Q\bar{Q}$ case

Additional final-state radiation  $\longrightarrow$  Massive emitters: only new **soft** divergencies

- $q_T$  subtraction successfully applied at NLO for  $t\bar{t}$  (and NNLO off-diagonal)

[Bonciani, Catani, Grazzini, Sargsyan, Torre, '15]

- At NNLO:

$$d\sigma_{\text{NNLO}}^{t\bar{t}} = \mathcal{H}_{\text{NNLO}}^{t\bar{t}} \otimes d\sigma_{\text{LO}}^{t\bar{t}} + \left[ d\sigma_{\text{NLO}}^{t\bar{t}+\text{jet}} - d\sigma_{\text{NNLO}}^{t\bar{t}, \text{CT}} \right]$$

Two-loop virtuals  
known numerically

[Baernreuther, Czakon, Fiedler, '13]

Structure of hard-virtual  
function was unknown!

Known from studies  
on  $q_T$  resummation

[Li, Li, Shao, Yang, Zu, '13;  
Catani, Grazzini, Torre, '14]

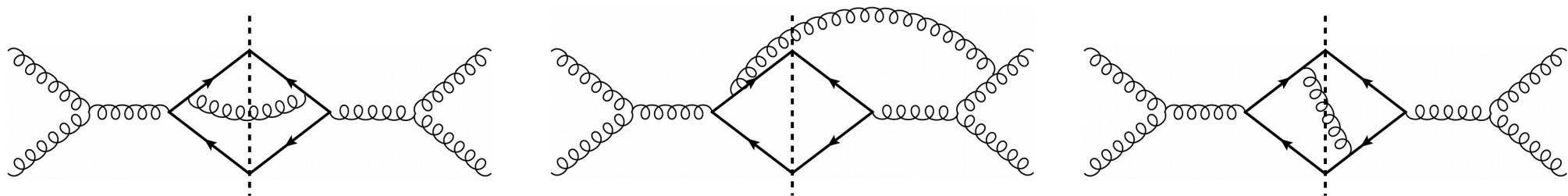
$$H \sim \langle \tilde{\mathcal{M}} | \tilde{\mathcal{M}} \rangle \longrightarrow (\mathbf{H}\Delta) \sim \langle \tilde{\mathcal{M}} | \Delta | \tilde{\mathcal{M}} \rangle$$

$$\left( \begin{array}{c} \text{Equivalent to:} \\ \tilde{I} \longrightarrow \tilde{\mathbf{I}} \end{array} \right)$$

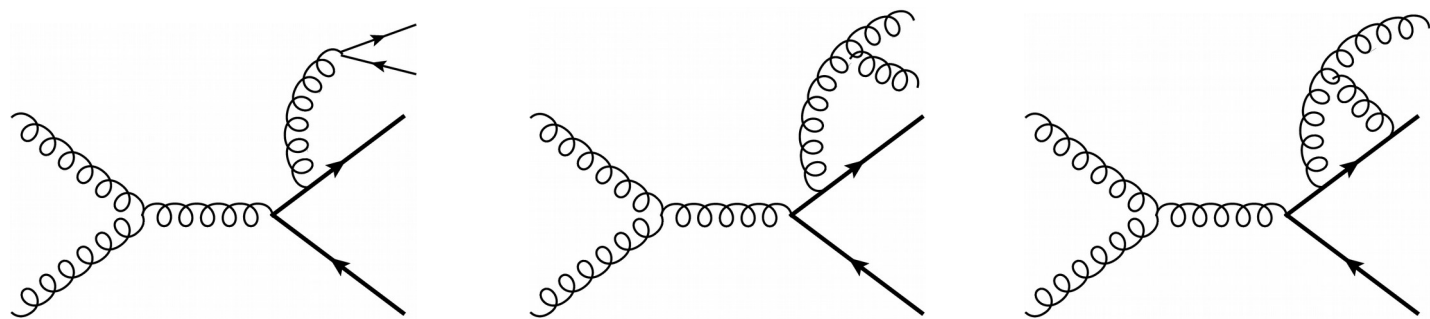
- Additional radiative soft factor  $\Delta$
- **NEW**: colour correlations in  $\Delta$  (4 hard partons) absent in the colourless case
- Its computation: integrate emissions from final state (specifically,  $d\sigma/d^2q_T$ )

Computation of  $d\sigma/d^2q_T$  at low  $q_T$  implies the PS integration of (appropriately defined) soft currents

At NLO we have single gluon emission contributions



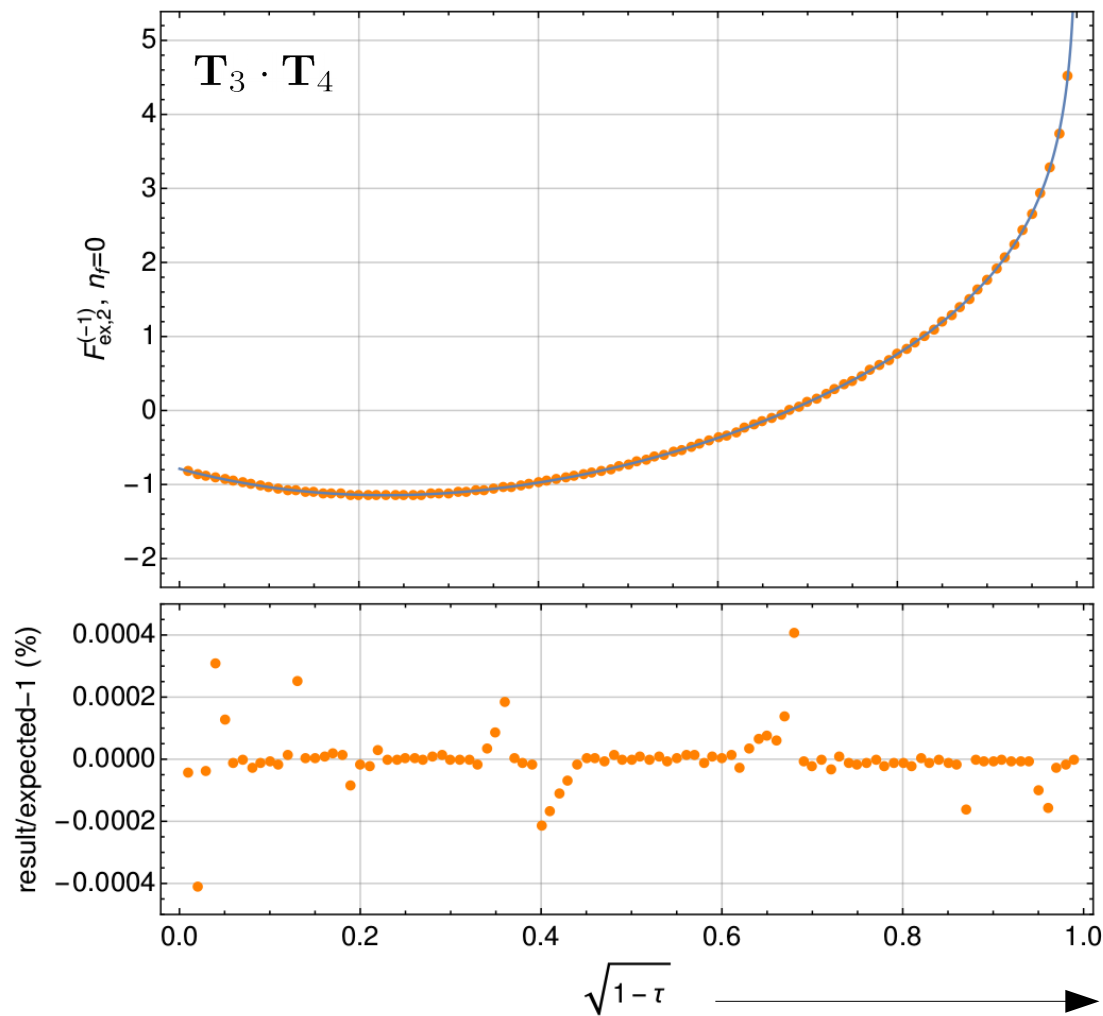
At NNLO we have  $q\bar{q}$ , double gluon and one-loop single gluon soft emissions



Most of the calculation was performed analitically, (numerically-)small piece obtained numerically (evaluated in 10000 PS points optimized for ttbar production)

# Final result - pole cancellation

- Final result:  $gg + qq + g$  (1-loop)  $\rightarrow$  triple poles cancel here
- Poles can be predicted, they cancel IR singularities of (a piece of) the 2-loop virtuals
- Analytic cancellation for all contributions, except for  $\mathbf{T}_3 \cdot \mathbf{T}_4$  single pole:



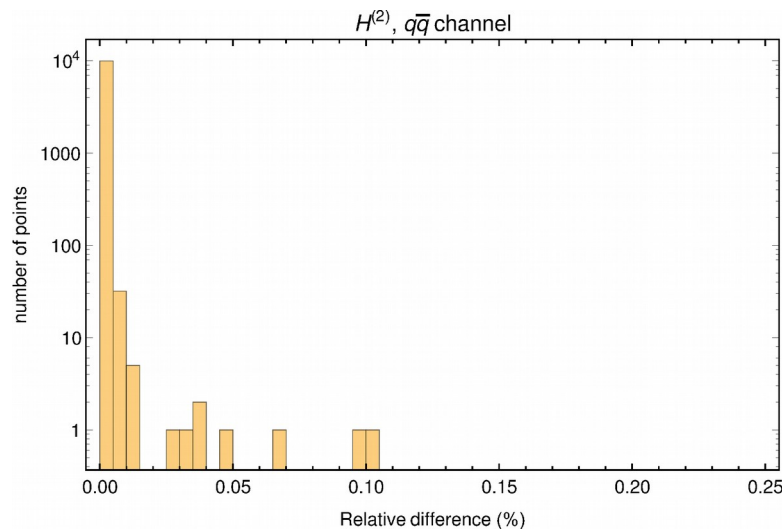
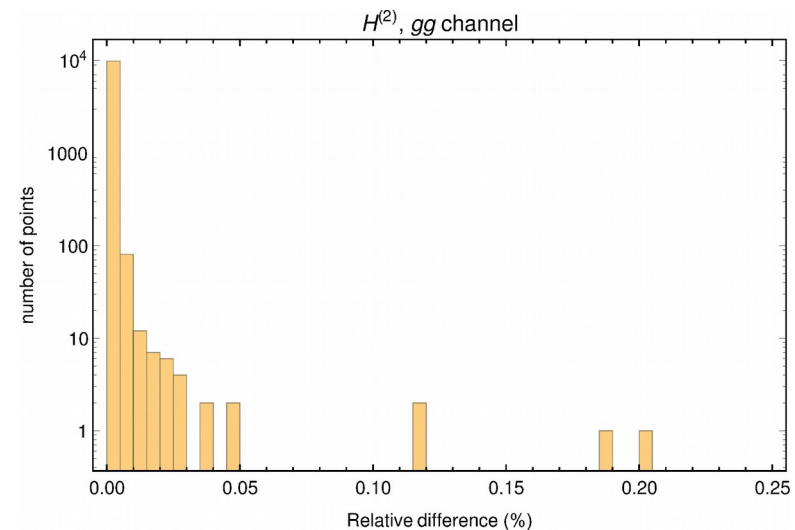
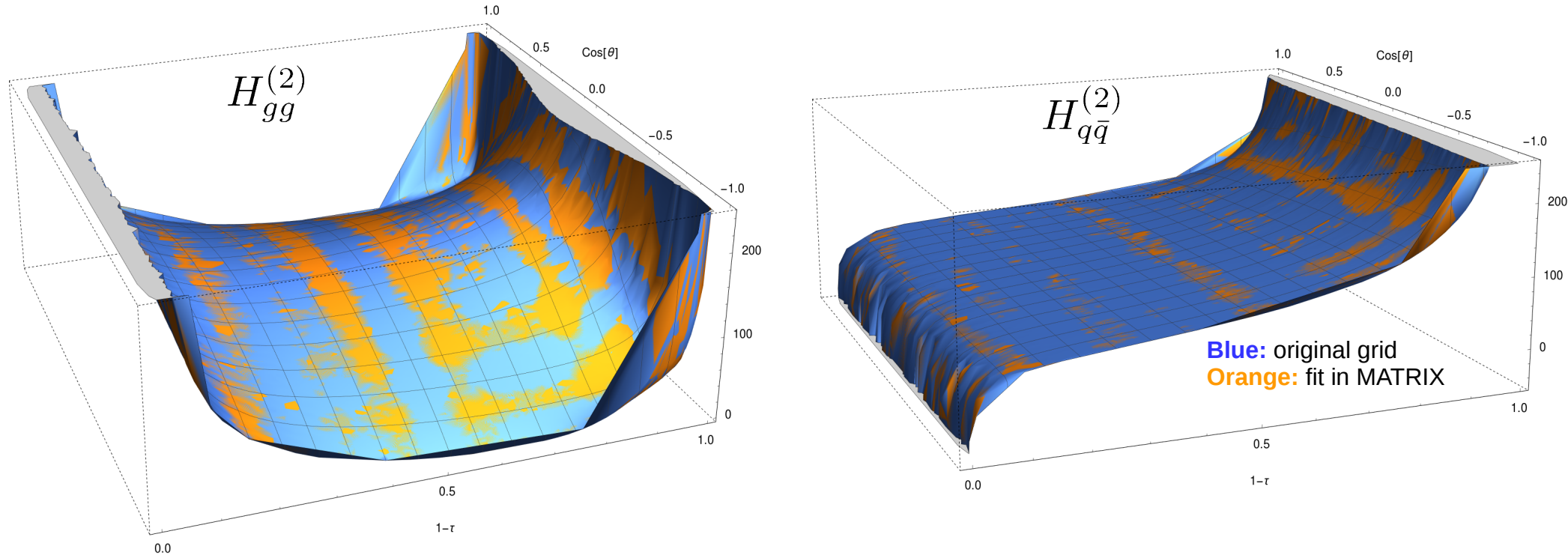
→ Poles independent of  $\cos(\theta)$



# Final result - $H^{(2)}$

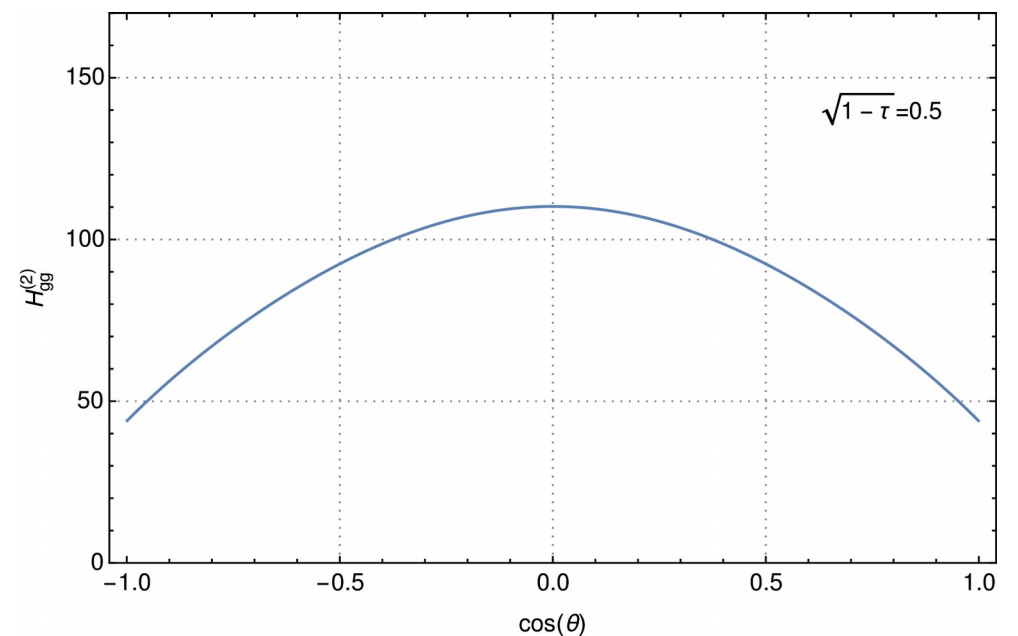
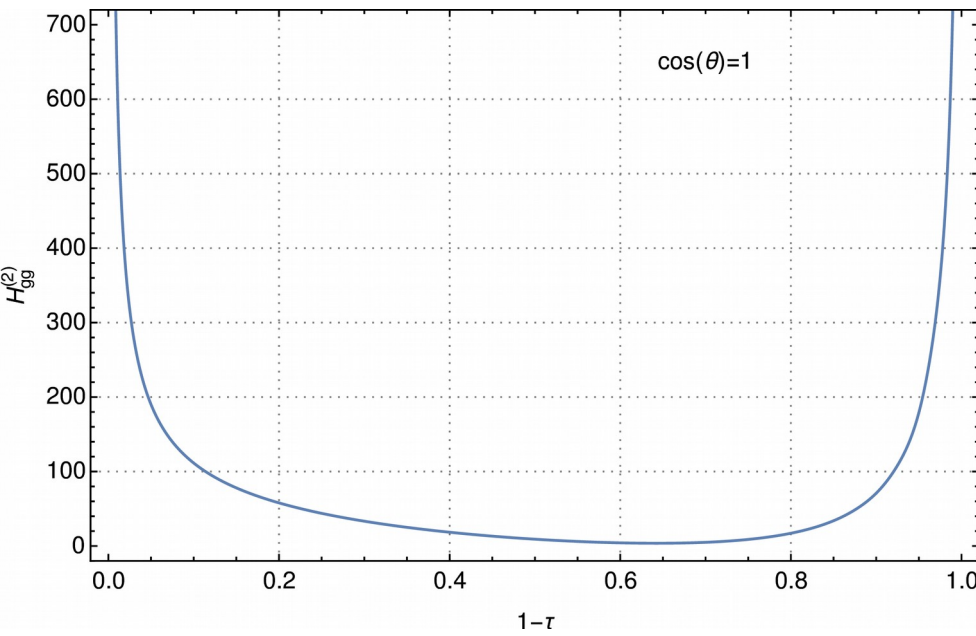
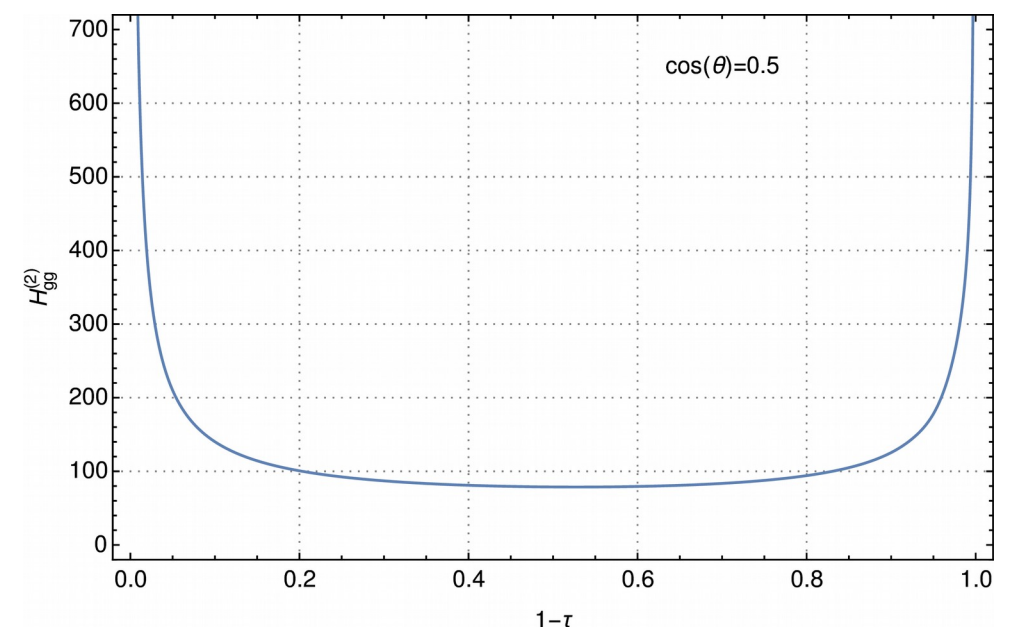
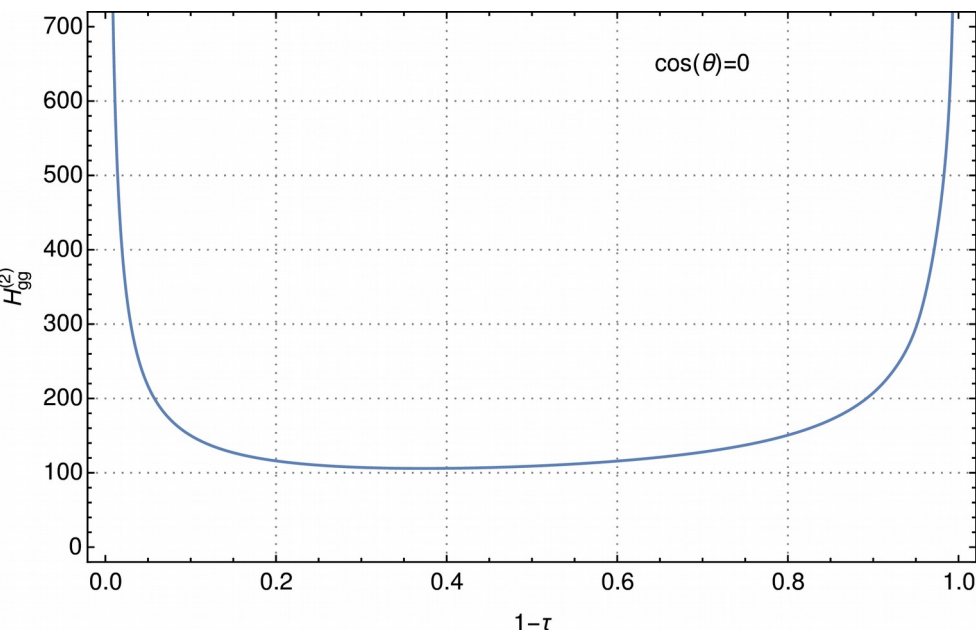
- Implementation in MATRIX: combination of analytic results + grids
- Final result for hard-virtual function  $H^{(2)}$ , including two-loop virtuals from [Baernreuther, Czakon, Fiedler, '13]

Small numerical impact:  
~3% of the NNLO piece at 13TeV



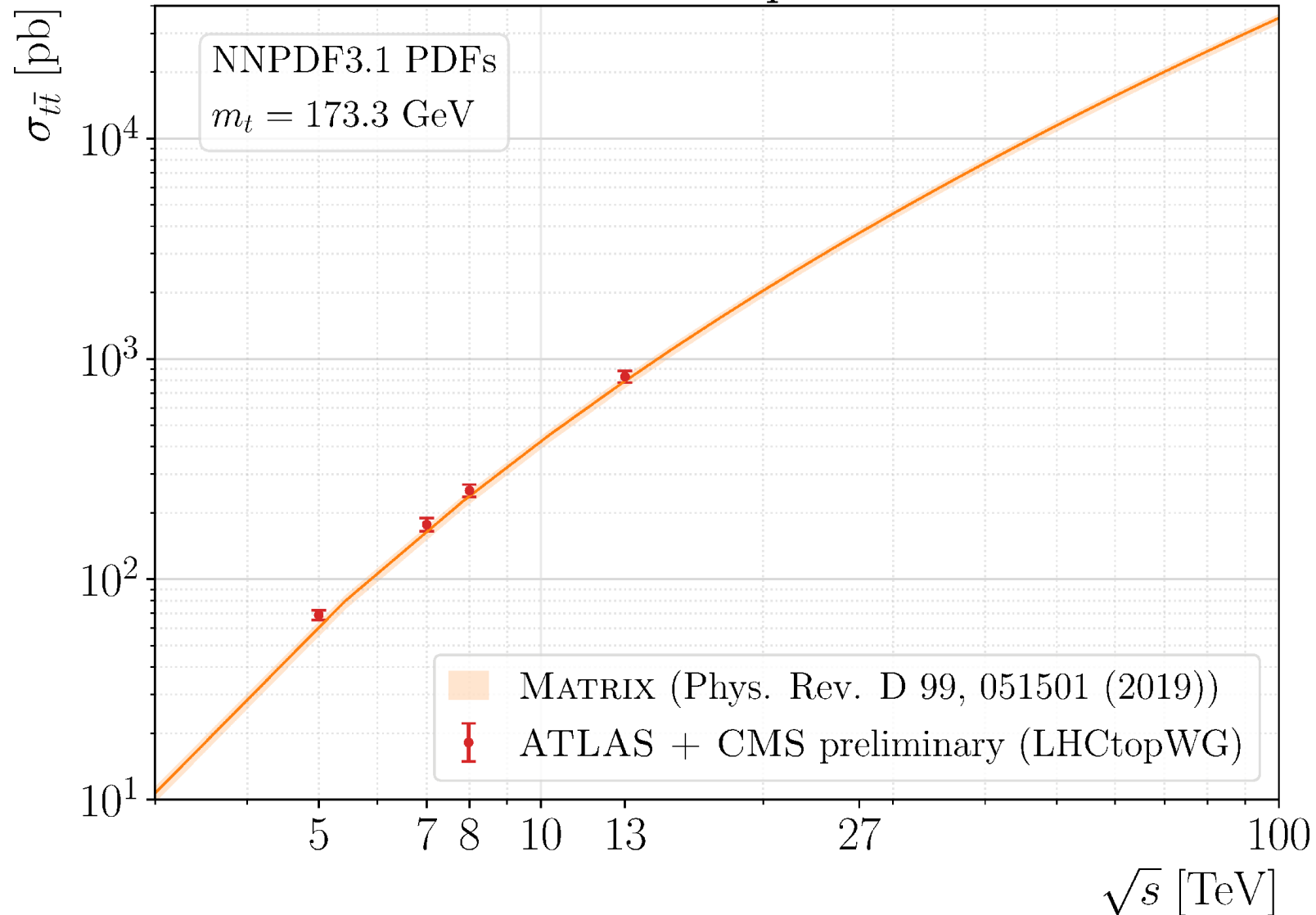
- Comparison of grid against fit shows it is well under control
- We also performed validations removing a large fraction of the grid points
- No need to include additional uncertainty due to numerical piece

# Final result - $H^{(2)}$



# Numerical results

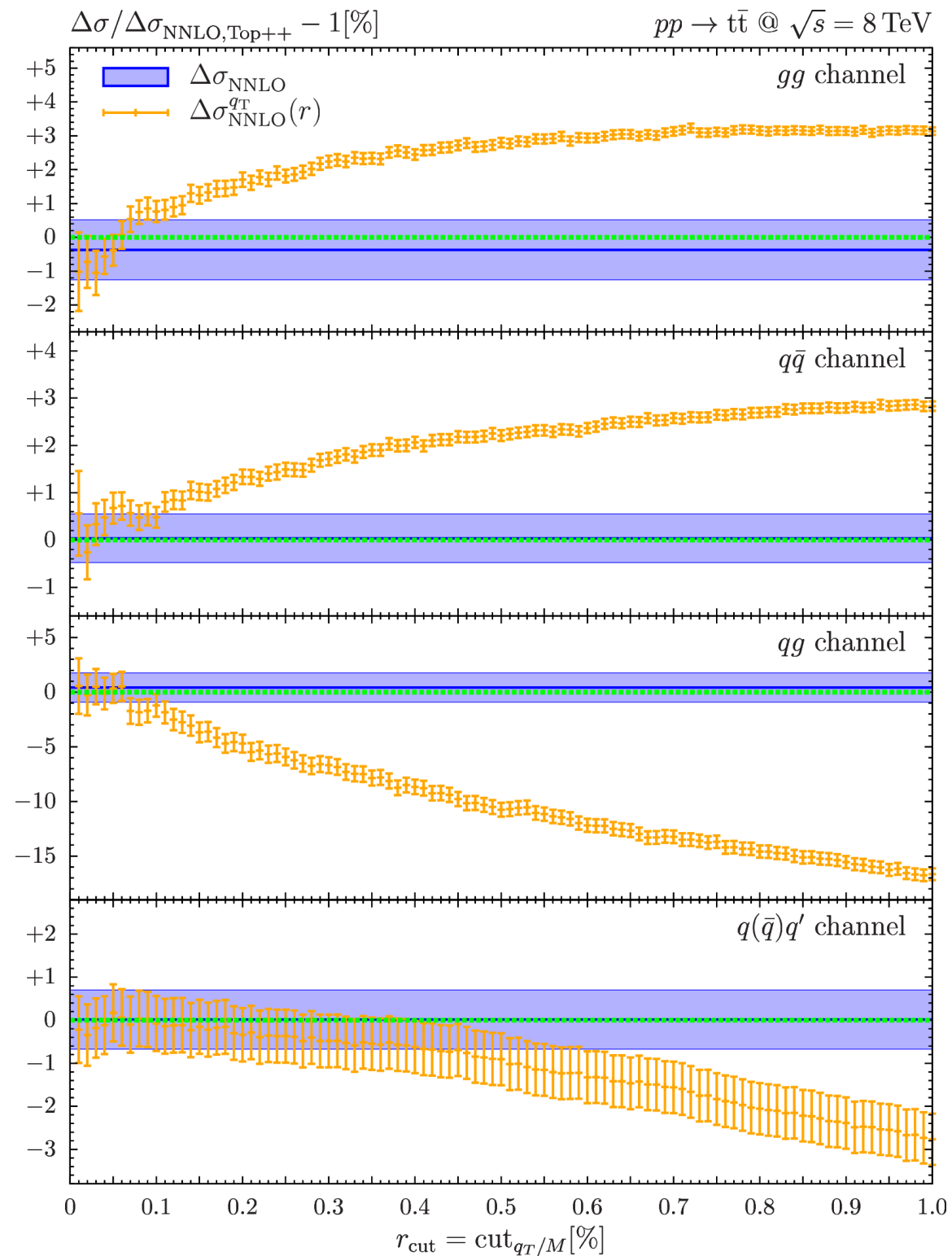
NNLO  $t\bar{t}$  cross section computed with MATRIX



Per-mille level accuracy for total cross section achievable in ~1000CPU days

# Numerical results

- Quality of the  $r_{\text{cut}} \rightarrow 0$  extrapolation can be studied looking at the  $r_{\text{cut}}$  dependence
- Stronger  $r_{\text{cut}}$  dependence compared to colourless case due to soft radiation (obs: not in qq')
- Extrapolation performed doing a quadratic least-squares fit
- Results for total cross section in full agreement with Top++



# Summary and outlook

- We performed a new calculation of the NNLO  $t\bar{t}$  cross section
- First complete application of  $q_T$  subtraction for  $Q\bar{Q}$  at NNLO

*Near and far future:*

- Complete the validation of fully differential results
- Provide a public tool (i.e. new MATRIX release) to compute them
- Improve NNLO QCD with other corrections: NLO EW, Coulomb resummation
- Include decays
- Perform NNLL  $q_T$  resummation (some ingredients still missing)
- Extend the calculation to  $t\bar{t}$ +colourless
  - Some contributions can be recycled (not the most difficult ones...)
  - First step: do it at NLO [Fabre, Grazzini, in progress]